

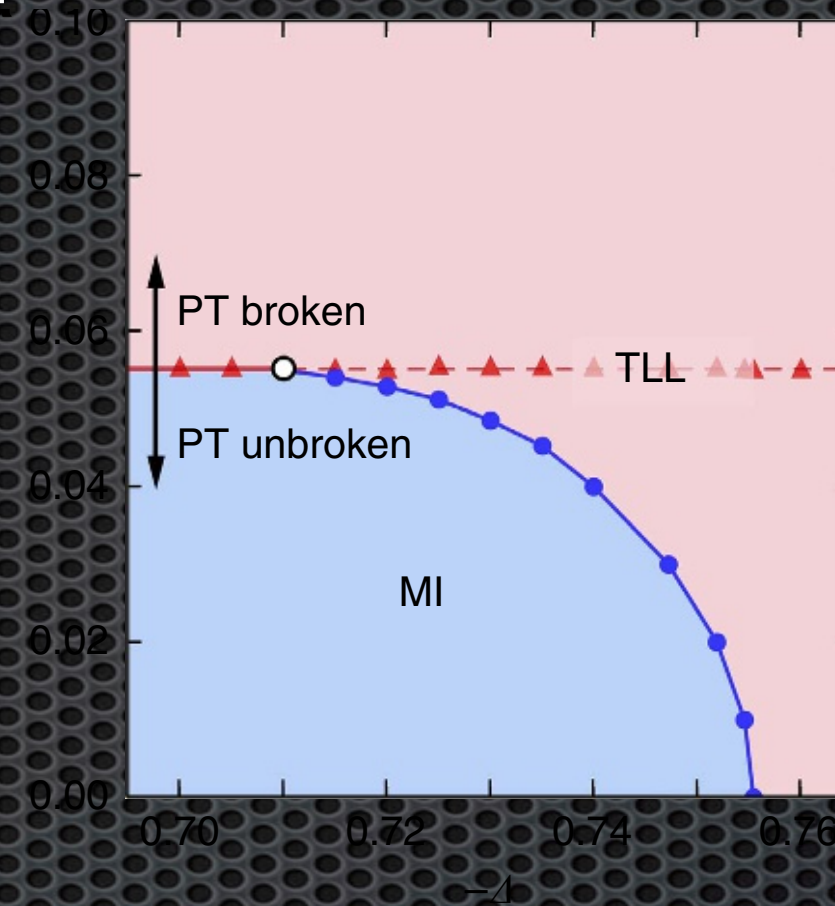
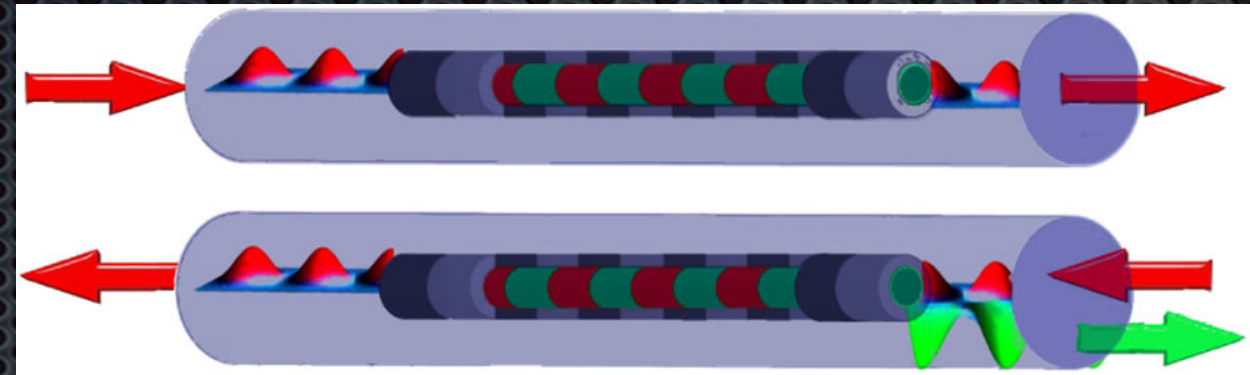
# Non-Hermitian Holography

Daniel Areán, Xixón, April 2023

Based on 1912.06647 w/ K. Landsteiner and I. Salazar; and in progress w/ D. García Fariña

# Non-Hermitian Holography. Outlook

- Non-Hermitian PT-symmetric QM

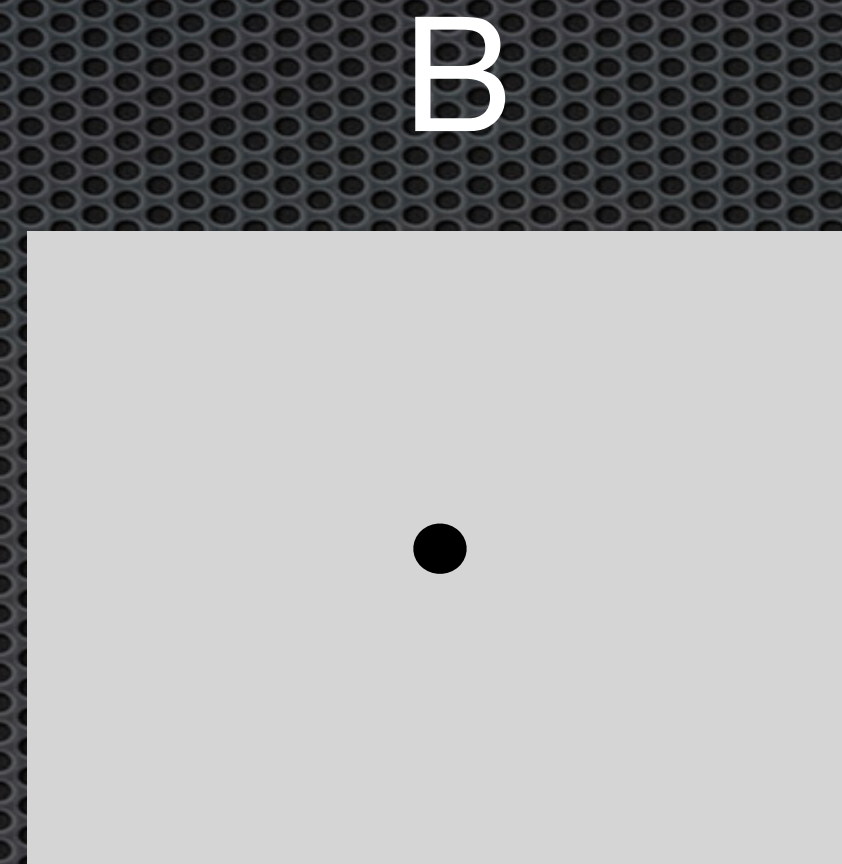
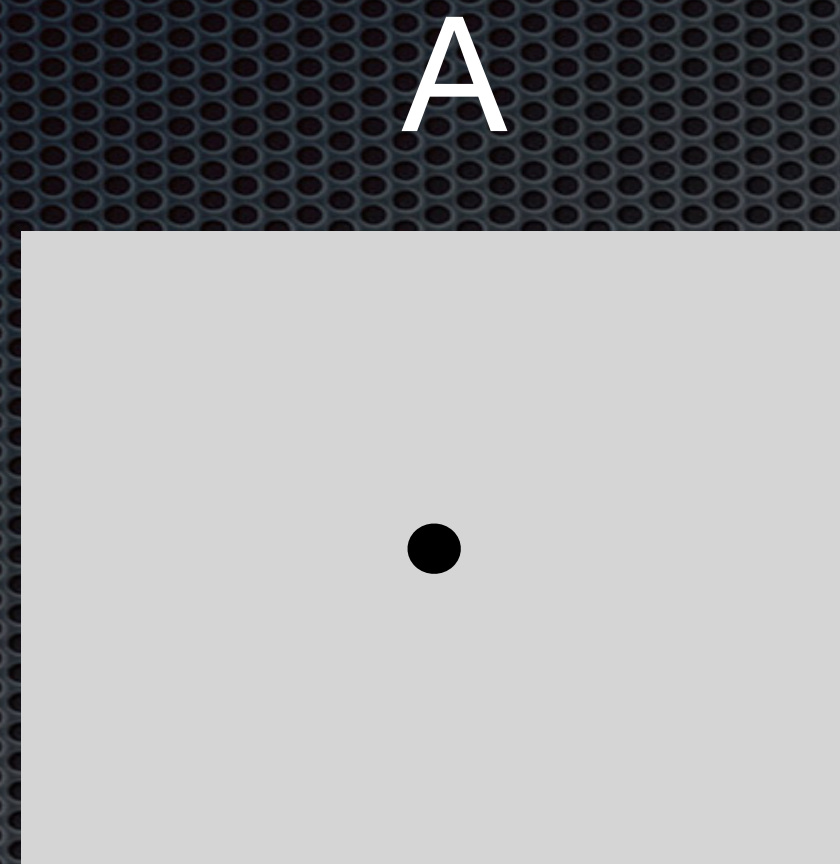


applications in optics, many body systems (cold atoms), quantum simulation, ...

- non-Hermitian PT-symmetric holography
- Modulated (non-Hermitian) configurations

# PT-symmetric QM. 2 state Hamiltonian

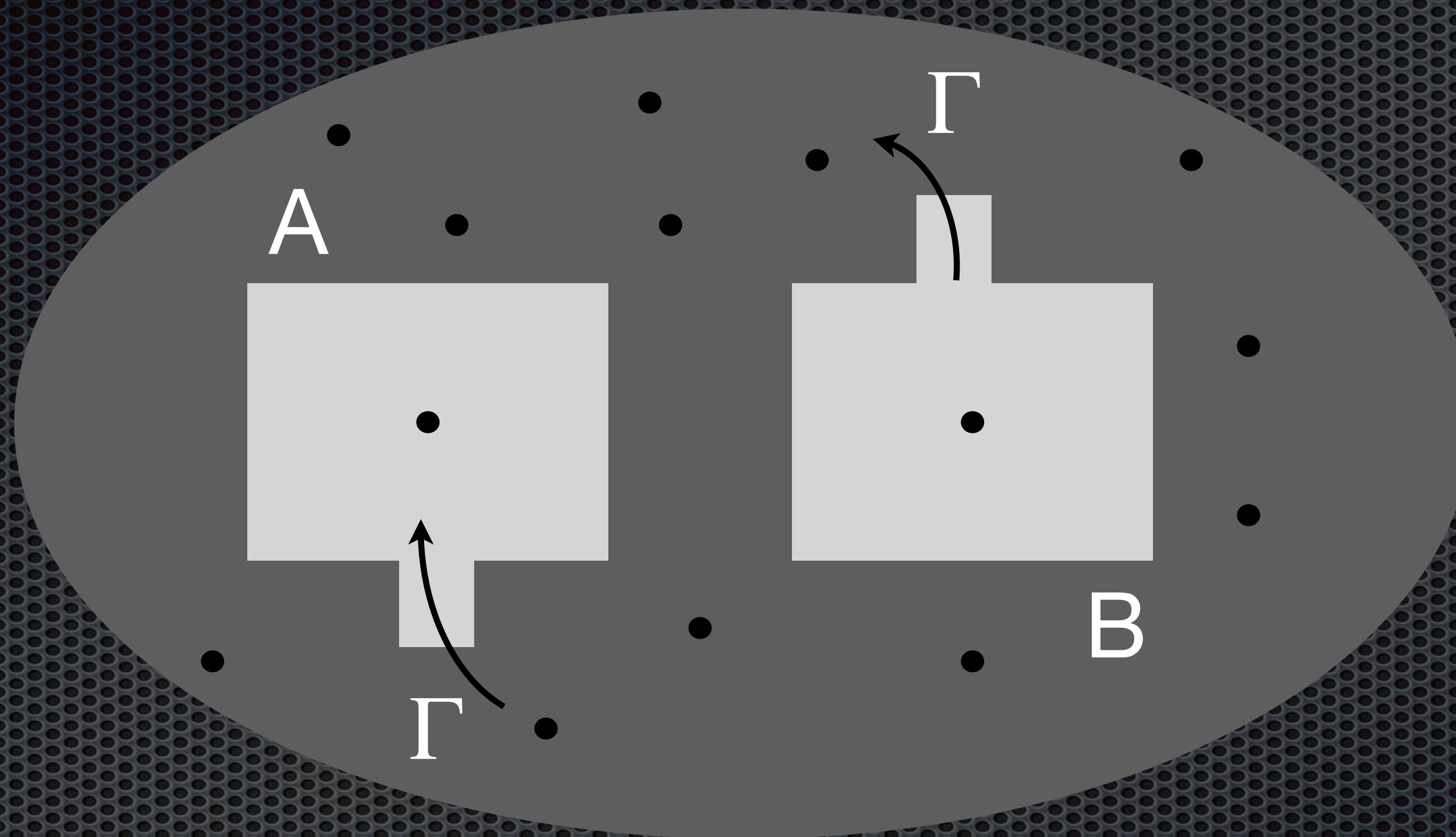
[Bender'97]



$$H = \begin{pmatrix} E & 0 \\ 0 & E \end{pmatrix}$$

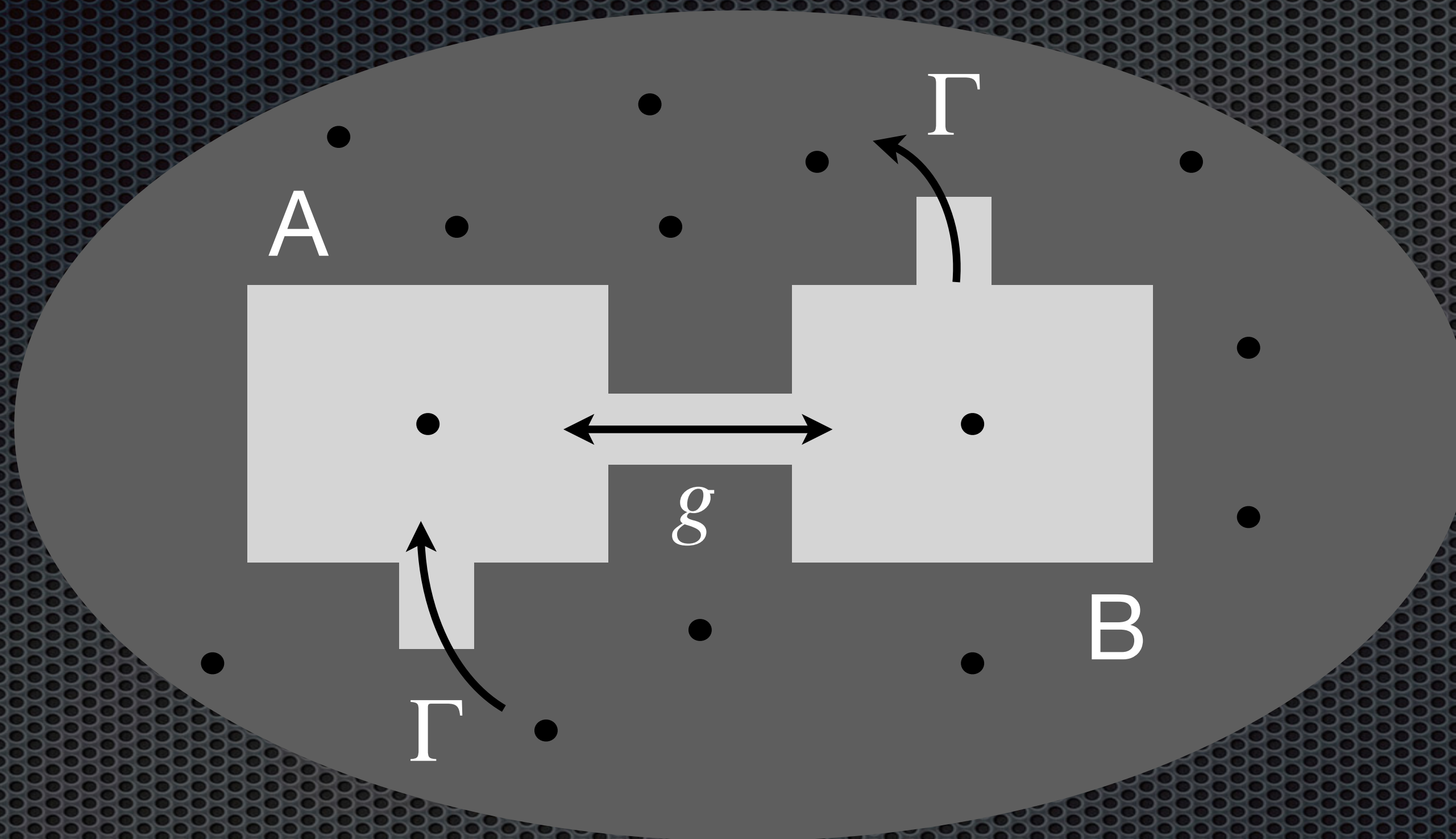
$$\psi_{A,B} = e^{-iEt}$$

# PT-symmetric QM. 2 state Hamiltonian



$$H = \begin{pmatrix} E - i\Gamma & 0 \\ 0 & E + i\Gamma \end{pmatrix} \quad \psi_{A,B} = e^{-(iE \pm \Gamma)t}$$

# PT-symmetric QM. 2 state Hamiltonian



$$H = \begin{pmatrix} E - i\Gamma & g \\ g & E + i\Gamma \end{pmatrix}$$

Eigenvalues

$$\lambda_{\pm} = E \pm \sqrt{g^2 - \Gamma^2}$$

# PT-symmetric QM

$$H = \begin{pmatrix} E - i\Gamma & g \\ g & E + i\Gamma \end{pmatrix} \quad \lambda_{\pm} = E \pm \sqrt{g^2 - \Gamma^2}$$

PT-symmetric phase:  $|g| > \Gamma$

PT-broken phase:  $|g| < \Gamma$

Exceptional point:  $g = \Gamma$

# PT-symmetric QM

$$T(H) = \begin{pmatrix} E + i\Gamma & g \\ g & E - i\Gamma \end{pmatrix} \quad P = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$P T(H) P = H$$

PT-symmetric phase:  $|g| > \Gamma$

PT-broken phase:  $|g| < \Gamma$

# PT-symmetric QM. Pseudo-hermitian H

[Mostafazadeh '02, '03, '20; Fring '22]

$$[H, PT] = 0, \quad PT|\psi\rangle = e^{i\phi}|\psi\rangle, \quad PT\lambda|\psi\rangle = \lambda^*|\psi\rangle$$

[ $\mapsto (PT)^2 = 1$ ]

H is pseudo-Hermitian (it has real eigenvalues)

Dyson map:  $\eta : H \longrightarrow h = \eta H \eta^{-1}, \quad h = h^\dagger$

$\mapsto$  metric  $\rho = \eta^\dagger \eta \Leftrightarrow \rho H^\dagger = H$

$$\langle \psi | \tilde{\psi} \rangle_\rho \equiv \langle \psi | \rho \tilde{\psi} \rangle \Rightarrow \langle \psi | H \tilde{\psi} \rangle_\rho = \langle H \psi | \tilde{\psi} \rangle_\rho$$



# PT-symmetric QM. Pseudo-hermitian H

$$H_2(\vec{g}) = E \mathbb{1} + \vec{g} \cdot \vec{\sigma} \iff H'_2(\vec{g}') = D(\vec{\alpha})^\dagger H_2(\vec{g}) D(\vec{\alpha})$$

$$SU(2) \text{ rotation } D(\vec{\alpha}) = e^{i\vec{\alpha} \cdot \vec{\sigma}/2} \implies \text{Dyson map: } \eta(\vec{\beta}) = D(-i\vec{\beta})$$

$$\text{E.g. } \vec{g} = (g', 0, 0), \vec{\alpha} = (0, i\beta, 0) \implies H_{\text{nH}} = E + g'(\cosh \beta \sigma_1 + i \sinh \beta \sigma_3)$$

$$\text{Eigenvalues } \lambda_{\pm} = E \pm g' \sqrt{\cosh^2 \beta - \sinh^2 \beta}$$

$$\text{Exceptional point } \begin{cases} \beta \rightarrow \infty \\ g' \rightarrow e^{-\beta} \tilde{g} \end{cases} \implies$$

# PT-symmetric QM. Dirac fermions

$$\text{Mass term: } \mathcal{L}_{\text{mass}} = M [\psi^\dagger \gamma^0 (1 - \gamma^5) \psi] + \bar{M} [\psi^\dagger \gamma^0 (1 + \gamma^5) \psi]$$

$$\mathcal{L}_{\text{mass}} = \mathcal{L}_{\text{mass}}^\dagger \leftrightarrow M = \bar{M}^*. \text{ Define } M = \frac{1}{2}(m_r + i m_i)$$

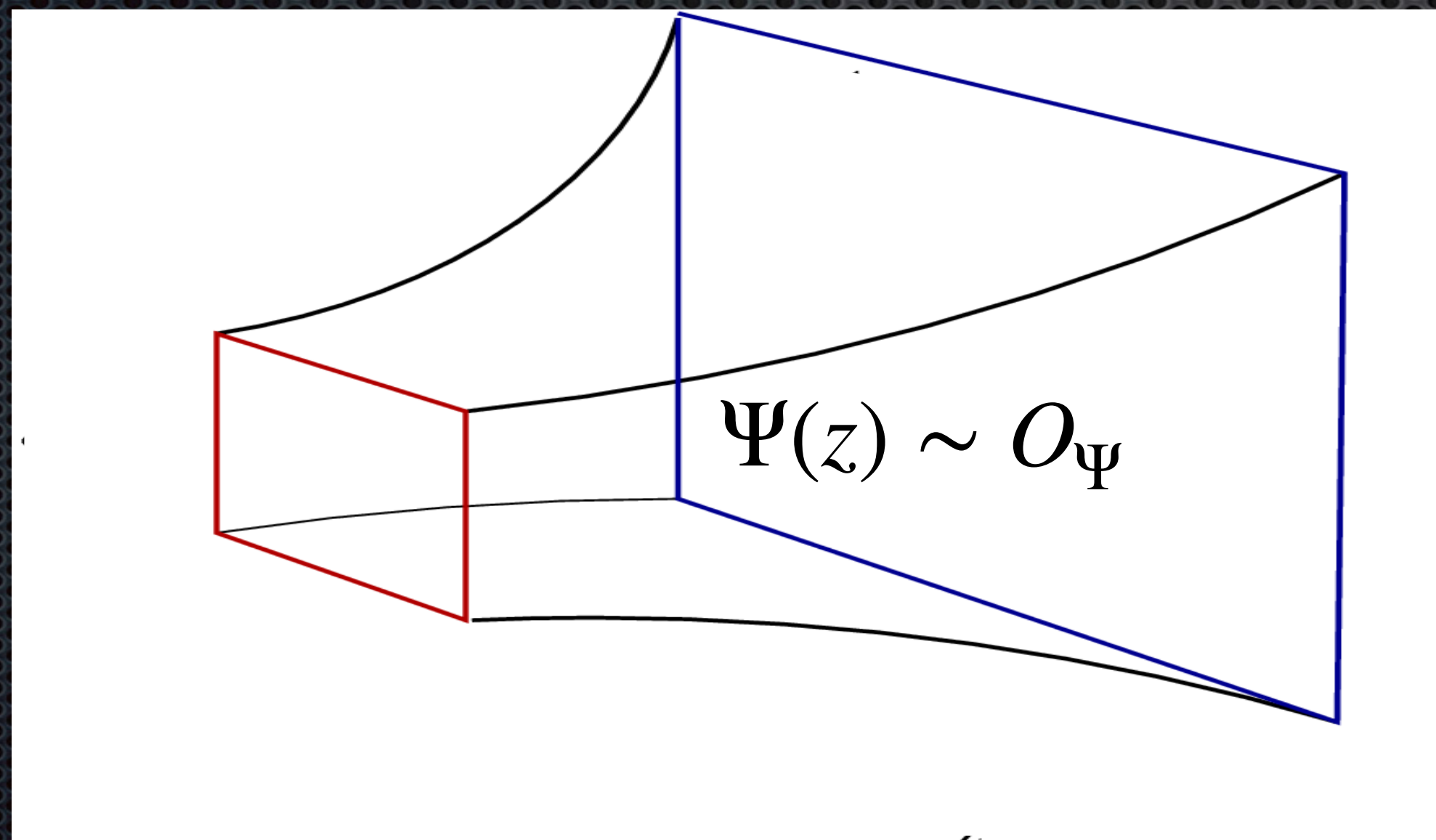
$$\text{Non-Hermitian 'rotation': } M \rightarrow e^{i\alpha} M, \bar{M} \rightarrow e^{-i\alpha} \bar{M} \text{ w/ } \alpha = i\beta$$

$$\hookrightarrow \mathcal{L}_{\text{mass}} \rightarrow \mathcal{L}_{\text{nH}} = m_r \cosh \beta \psi^\dagger \gamma^0 \psi - m_i \sinh \beta \psi^\dagger \gamma^0 \gamma^5 \psi$$

$$\text{Holography: complex operator } \mathcal{O} \sim \Psi(r) \text{ where } \Psi|_{\text{boundary}} = M$$

# Non-Hermitian Holography

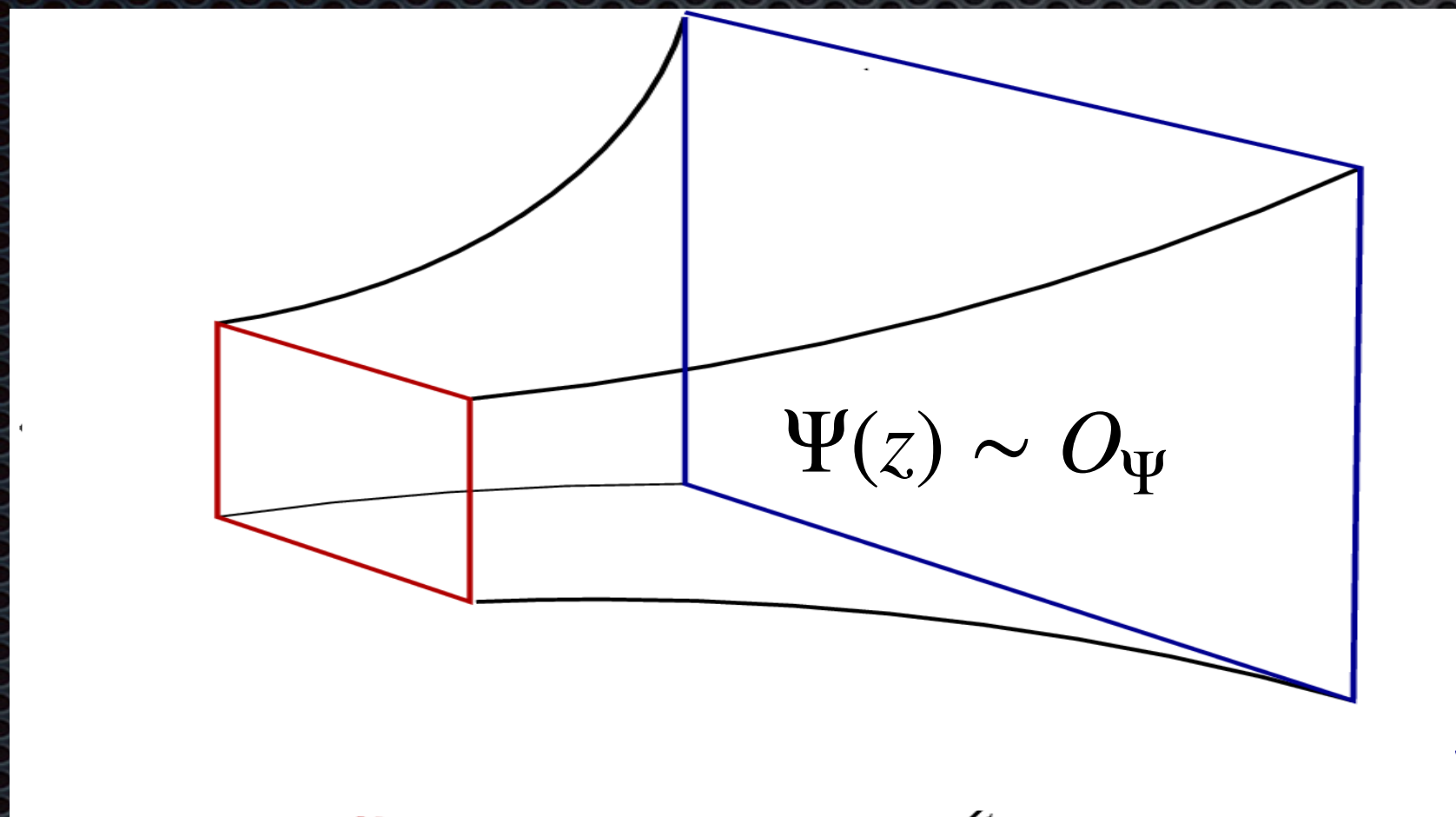
Complex scalar field  $\Psi$  in  $\sim$  AdS geometry:  $\Psi|_{\text{boundary}} = M$ ,



$$\sim \mathcal{L} = \dots + \bar{M} \langle O \rangle + M \langle O^\dagger \rangle$$

# Non-Hermitian Holography

Complex scalar field  $\Psi$  in  $\sim$  AdS



Dyson map:  $U(1)$  w/ imaginary phase

$$\Psi \sim M \rightarrow \Psi \sim M e^{-\theta}$$

$$\bar{\Psi} \sim M \rightarrow \Psi \sim M e^{\theta}$$

$$S = \int d^{d+1}x \sqrt{-g} \left[ R - 2\Lambda + F^2 + \frac{1}{2} g^{ab} (D_a \Psi \bar{D}_b \bar{\Psi} + a \leftrightarrow b) - m^2 \bar{\Psi} \Psi + v (\bar{\Psi} \Psi)^2 \right]$$

# Holography w/ non-Hermitian boundary conditions

$$S = \int d^{d+1}x \sqrt{-g} \left[ R - 2\Lambda + F^2 + \frac{1}{2} g^{ab} (D_a \Psi \bar{D}_b \bar{\Psi} + a \leftrightarrow b) - m^2 \bar{\Psi} \Psi + \nu (\bar{\Psi} \Psi)^2 \right]$$

UV ( $r \sim 0$ ) boundary conditions ( $m^2 = -2$ )

$$\Psi = M r + \nu r^2 + o(r^3)$$

$$\bar{\Psi} = M r + \nu r^2 + o(r^3)$$



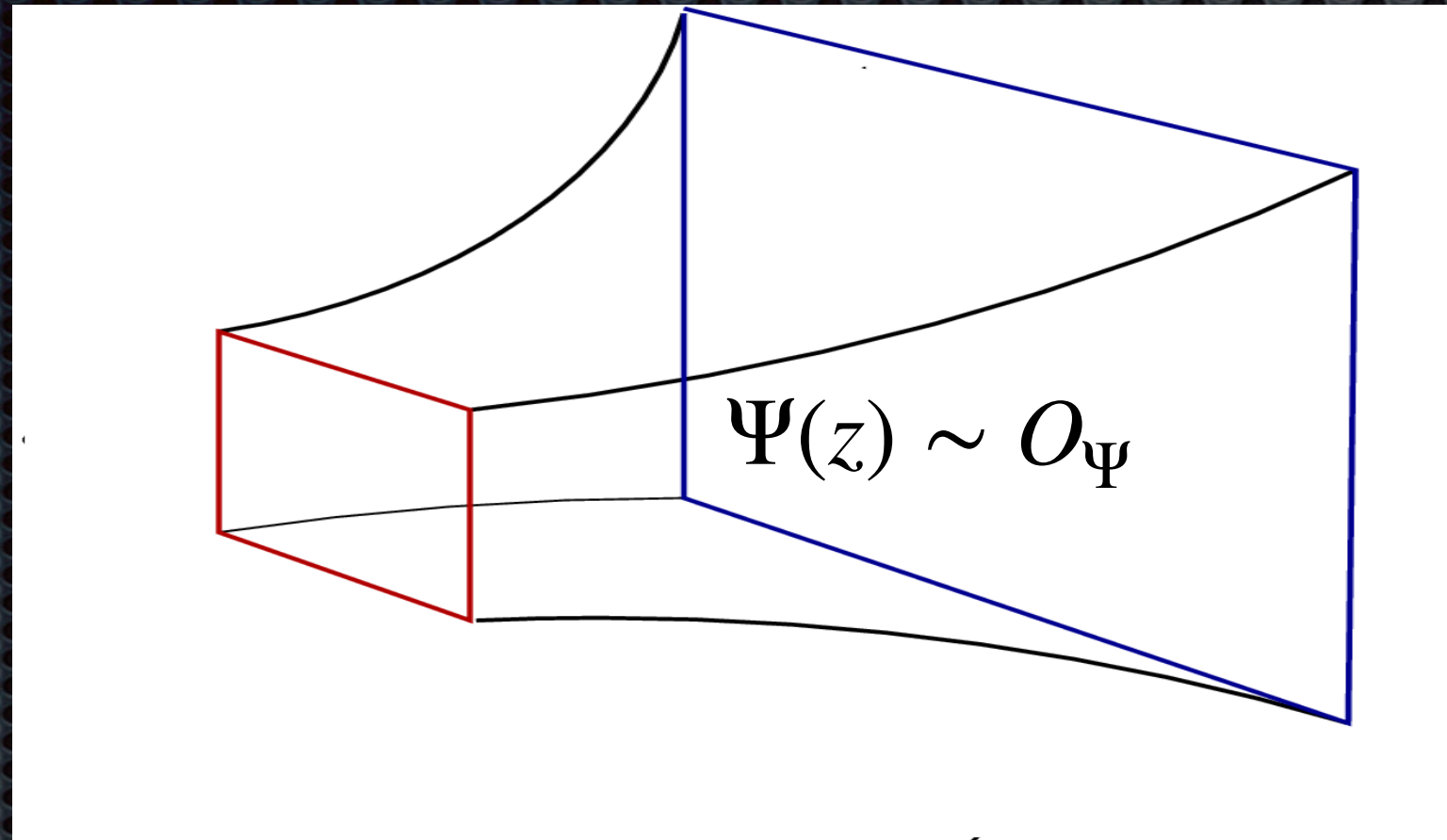
$$\Psi_{\text{boundary}} \rightarrow e^{i\alpha} \Psi_{\text{boundary}}$$

$$\bar{\Psi}_{\text{boundary}} \rightarrow e^{-i\alpha} \bar{\Psi}_{\text{boundary}}$$

Holographic Dyson map:  $\alpha = i\beta$

# Holography w/ non-Hermitian boundary conditions

UV ( $r \sim 0$ ) non-Hermitian boundary conditions ( $m^2 = -2$ )



$$\Psi = M e^{-\beta} r + \dots = (1 - \xi) s r + \dots$$

$$\bar{\Psi} = M e^{\beta} r + \dots = (1 + \xi) s r + \dots$$

Pseudo-Hermitian regimen:  $\xi < 1$

Exceptional point:  $\xi = 1$

non-Hermitian regime:  $\xi > 1$

$$e^{2\beta} = \frac{1 + \xi}{1 - \xi}$$

$$M = \sqrt{1 - \xi^2}$$

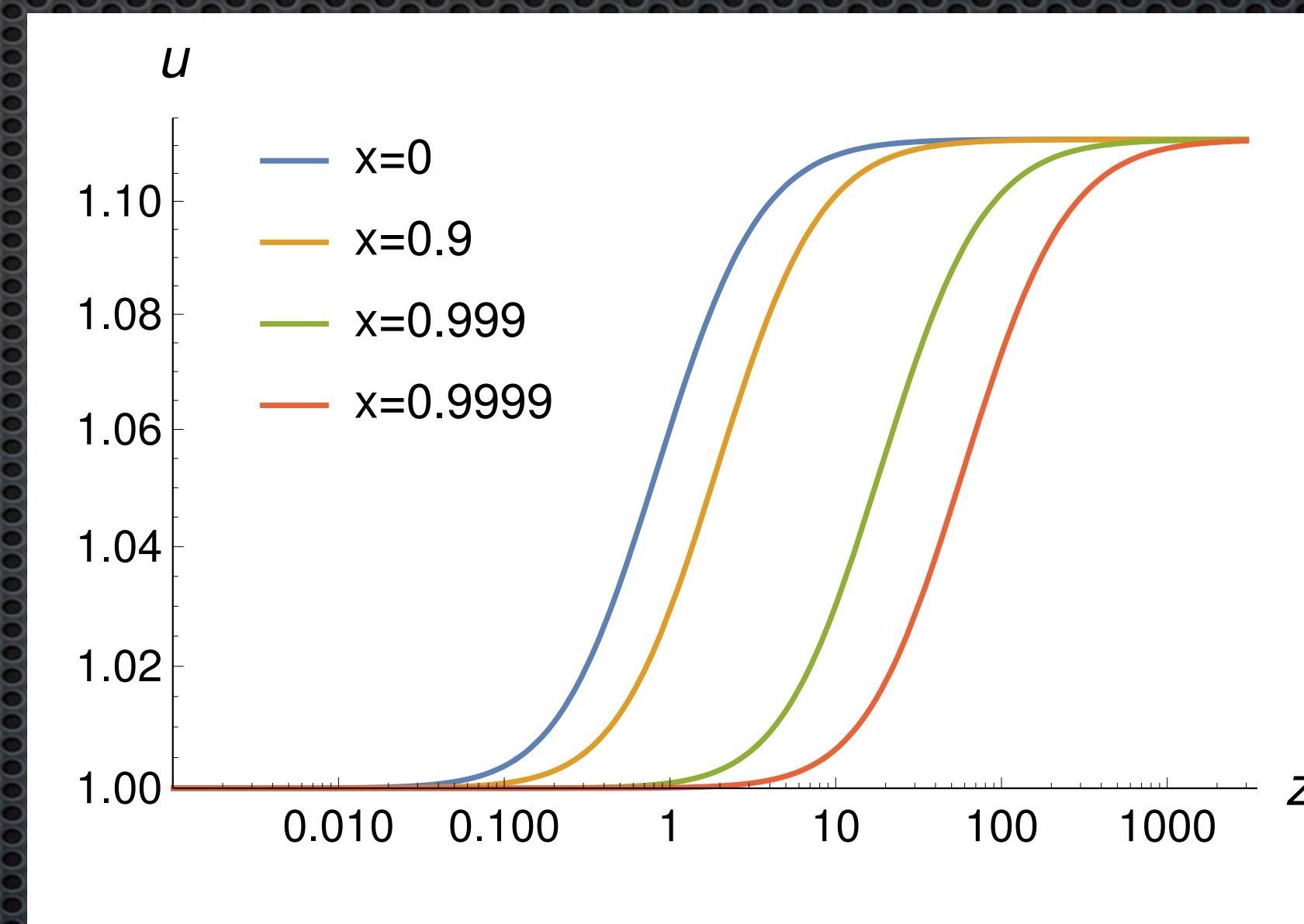
# Non-Hermitian Holography [1912.06647]

Pseudo-Hermitian domain walls:  $\xi < 1$

$T = 0$  AdS<sub>4</sub> to AdS<sub>4</sub> domain walls

$$\Psi \sim s(1 - \xi)r$$

$$\bar{\Psi} \sim s(1 + \xi)r$$



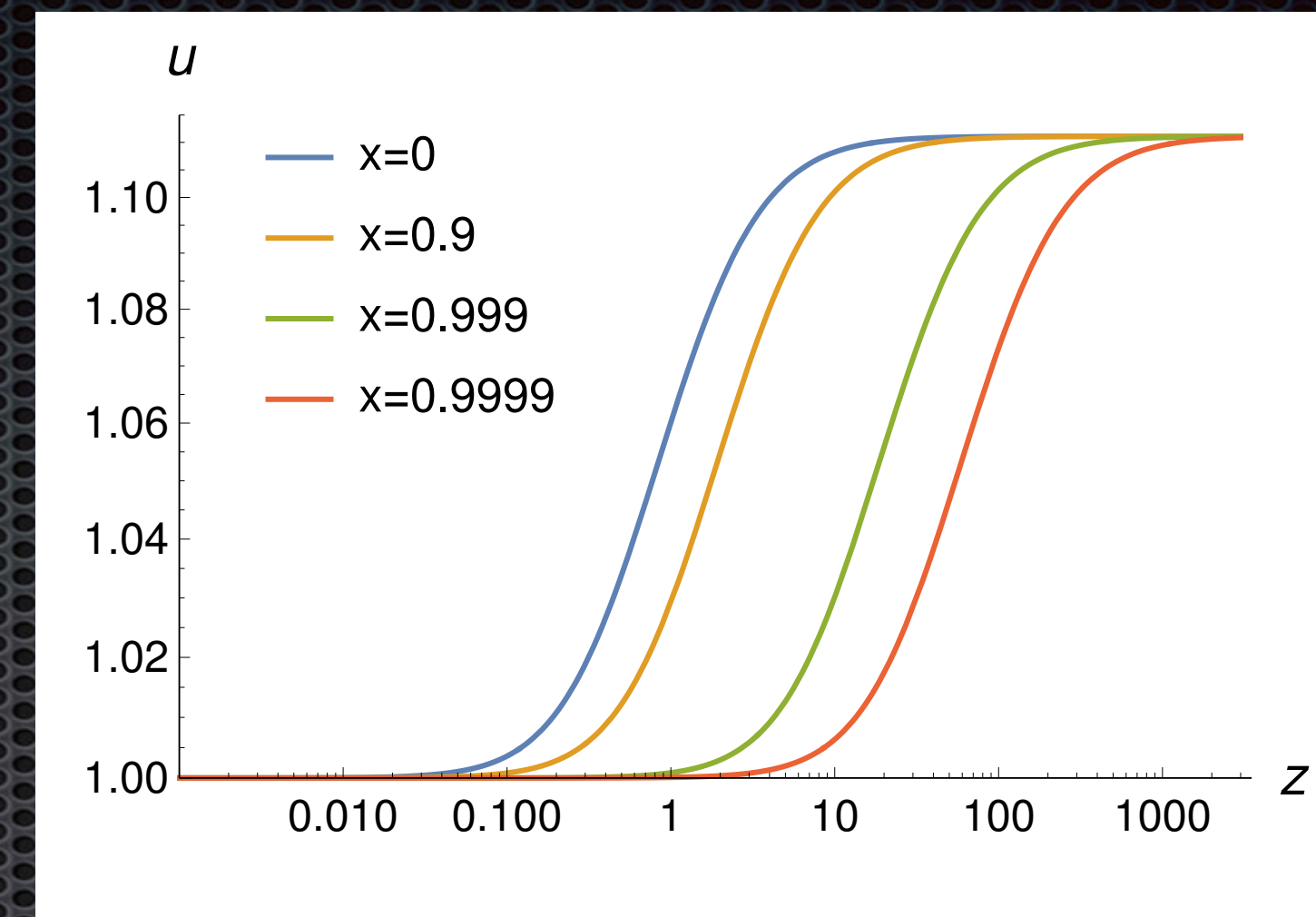
# Non-Hermitian Holography [1912.06647]

Holographic dyson map ( $\xi < 1$ )

Bulk action invariant under

$$\Psi(r) \rightarrow e^{\beta} \Psi(r)$$

$$\bar{\Psi}(r) \rightarrow e^{-\beta} \bar{\Psi}(r)$$



Pseudo-Hermitian

$$\Psi \sim s(1 - \xi) r$$

$$\bar{\Psi} \sim s(1 + \xi) r$$

$\sim$

Hermitian

$$\Psi \sim s \sqrt{1 - \xi^2} r$$

$$\bar{\Psi} \sim s \sqrt{1 - \xi^2} r$$

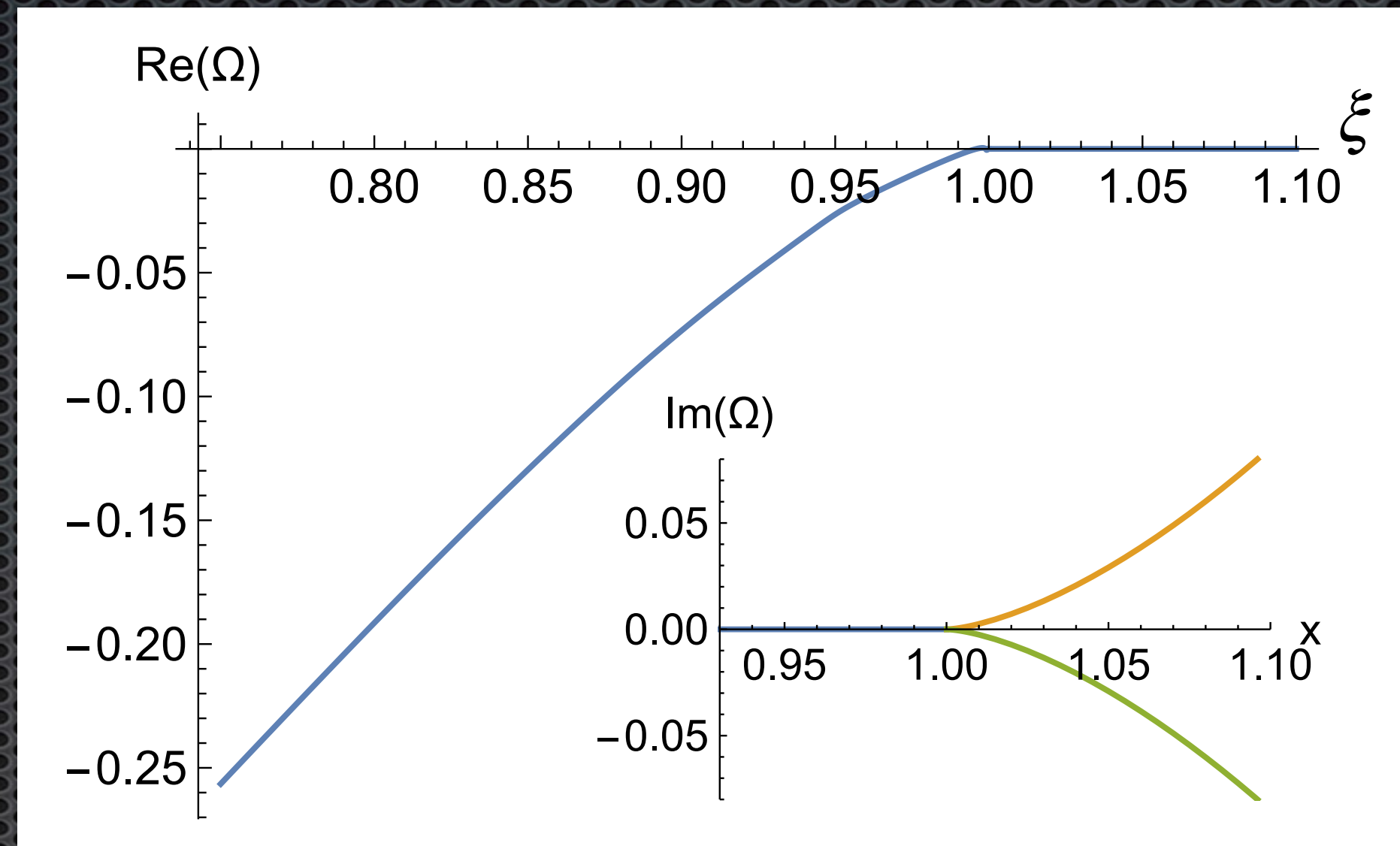


# Non-Hermitian Holography [1912.06647]

$$\Psi \sim s(1 - \xi)r$$

$$\bar{\Psi} \sim s(1 + \xi)r$$

Pseudo Hermitian for  $\xi < 1$



PT-symmetric phase:  $\xi < 1 \quad \sim \quad \xi = 0$

PT-broken phase:  $\xi > 1 \Rightarrow$  complex backgrounds

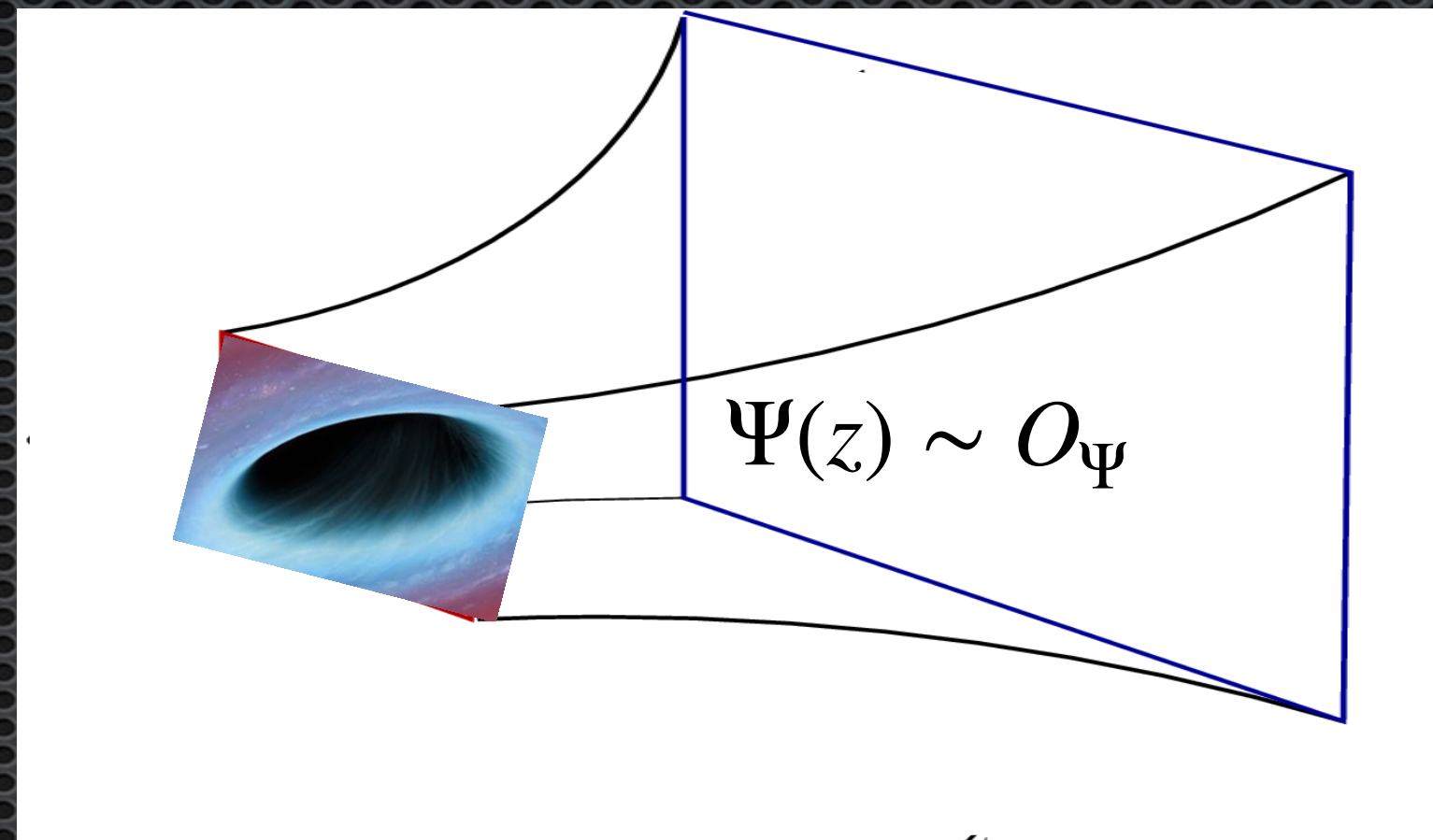
# Non-Hermitian Holography [1912.06647]

## Finite Temperature

Black brane geometry  $\implies T/s \neq 0$

$$\Psi \sim s(1 - \xi)r$$

$$\bar{\Psi} \sim s(1 + \xi)r$$

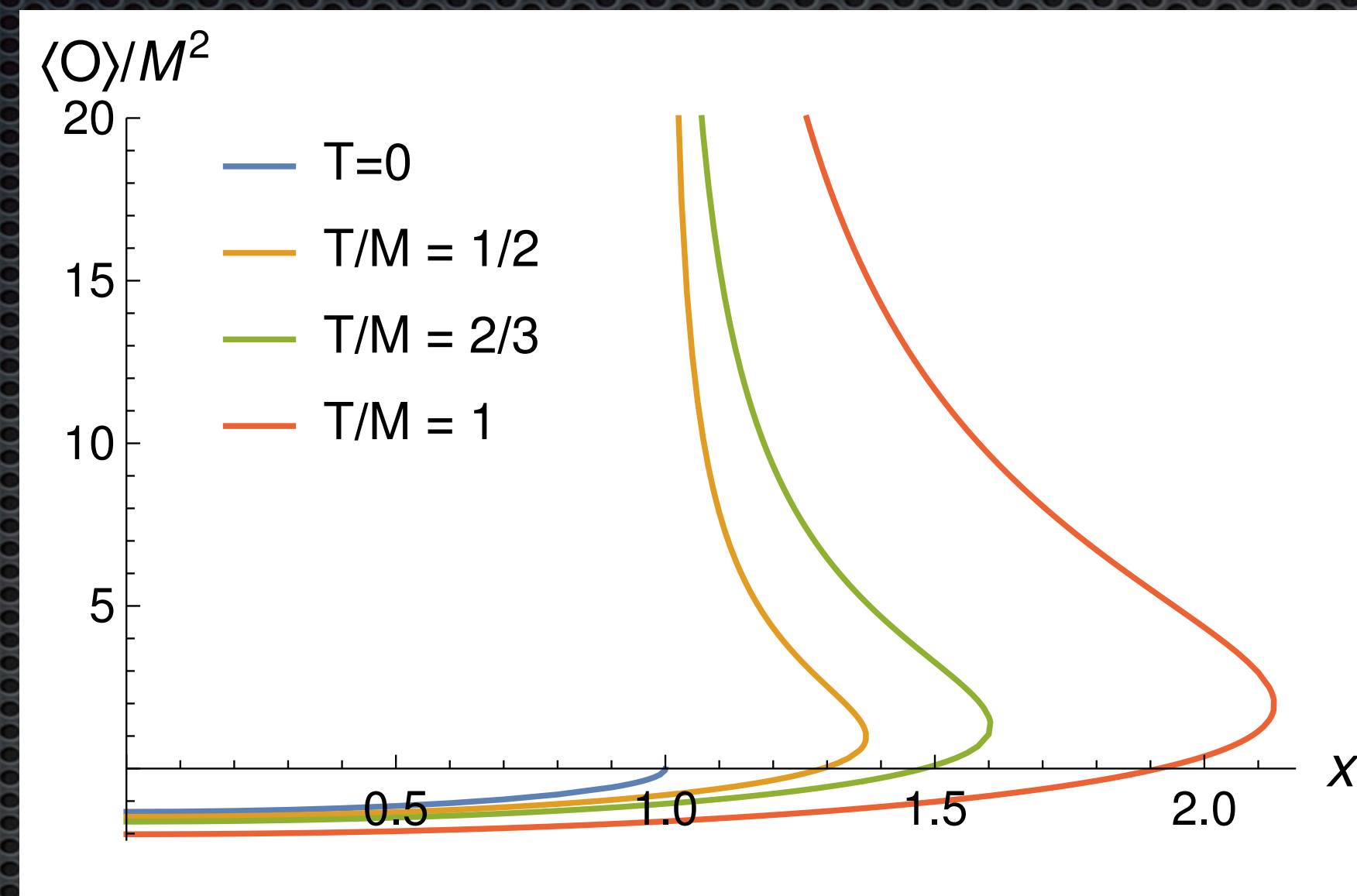


Real geometries up to  $x_c > 1$   
(unstable solutions)

# Non-Hermitian Holography [1912.06647] Finite Temperature

$$\Psi \sim s(1 - \xi)r$$

$$\bar{\Psi} \sim s(1 + \xi)r$$



Real geometries up to  $x_c > 1$

$$\xi < 1$$

Pseudo-Hermitian phase

$$\xi = 1$$

Exceptional point:  $\text{AdS}_4$

$$1 < \xi < \xi_c$$

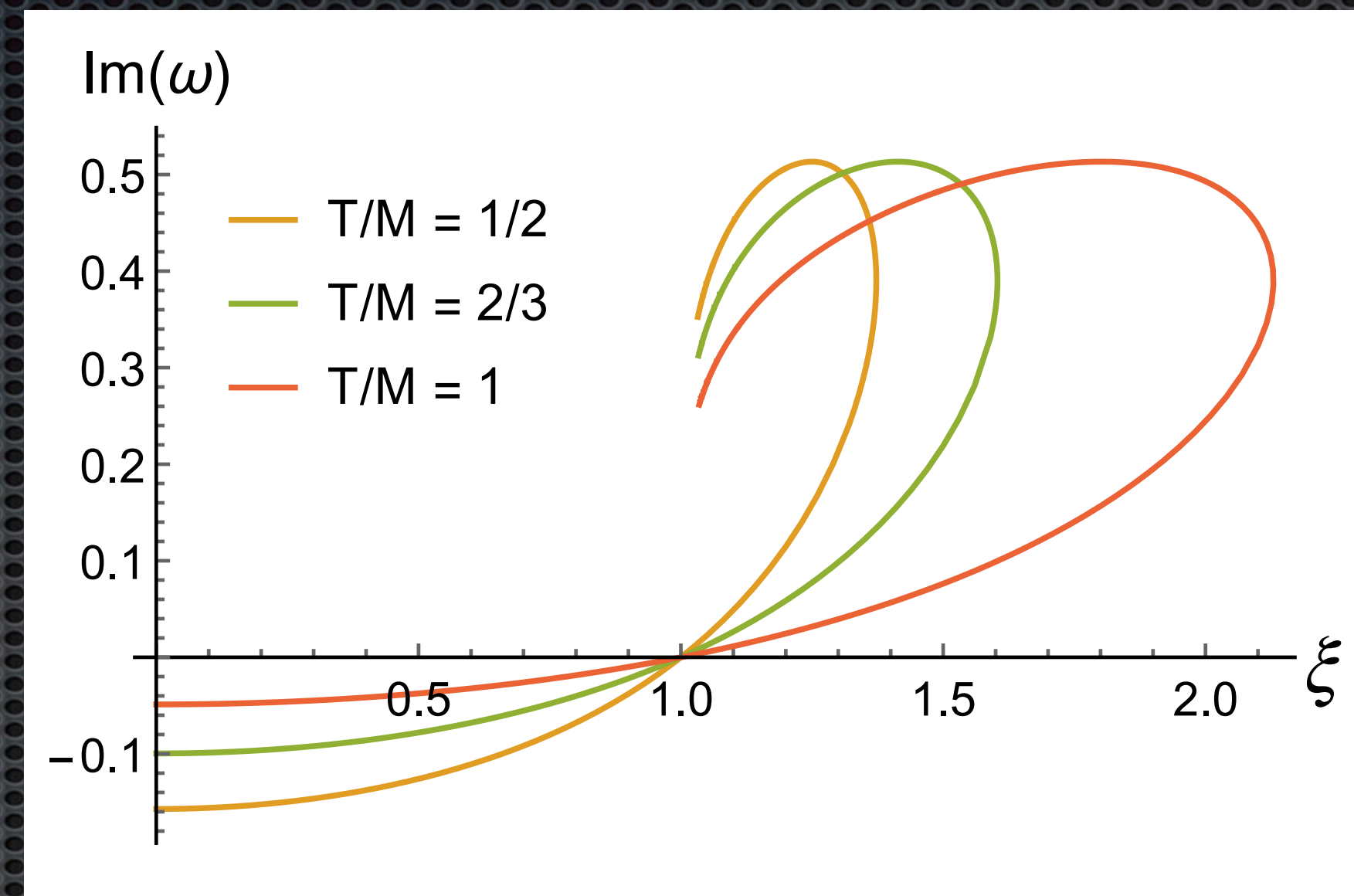
Two new (unstable) solutions

[see also Xian et al'23]

# Non-Hermitian Holography [1912.06647] Finite Temperature

$$\Psi \sim s(1 - \xi)r$$

$$\bar{\Psi} \sim s(1 + \xi)r$$



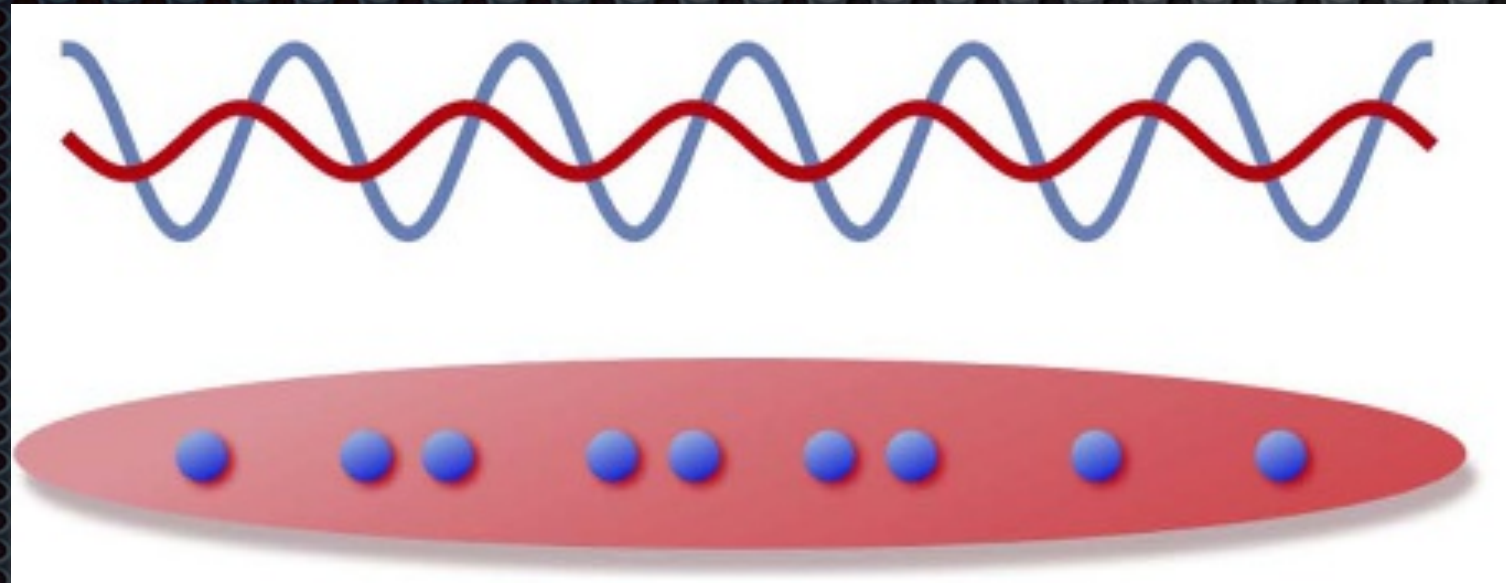
Real geometries up to  $x_c > 1$

$$1 < \xi < \xi_c \implies \text{Im}(\omega_{\text{QNM}}) > 0$$

# Non-Hermitian inhomogeneous holography

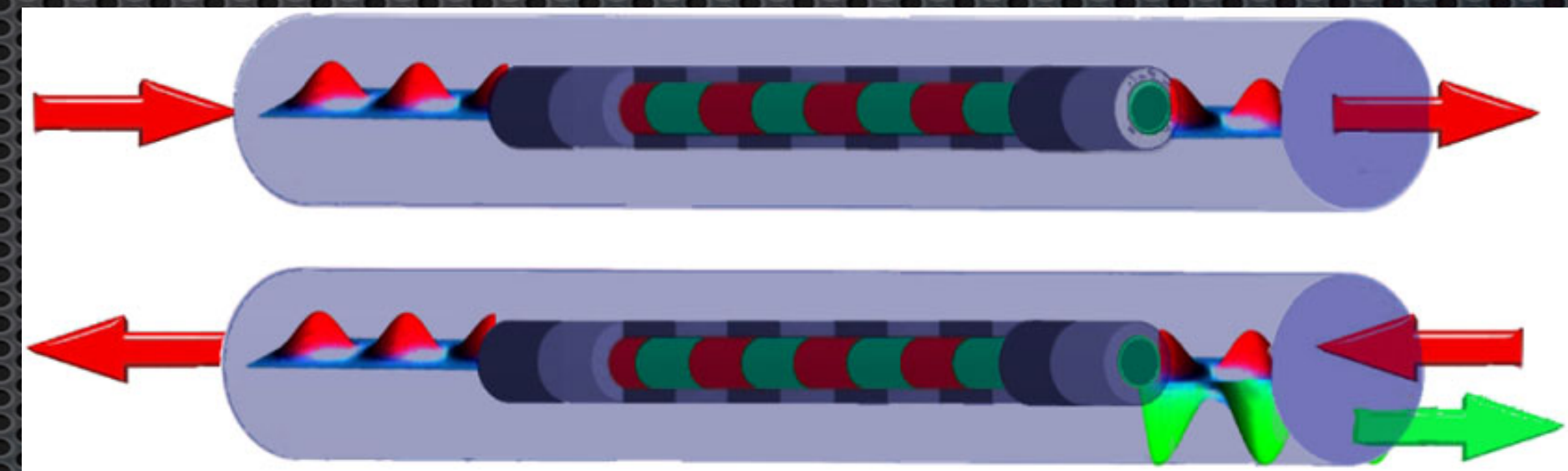
Space dependent non-Hermitian deformation

[Ashida et al '17]



New PT-breaking quantum critical points

[Lin et al '11]



PT-symmetric induced invisibility

➡ Non-Hermitian modulated BCs

$$\Psi \sim s(1 - \xi(x))r$$

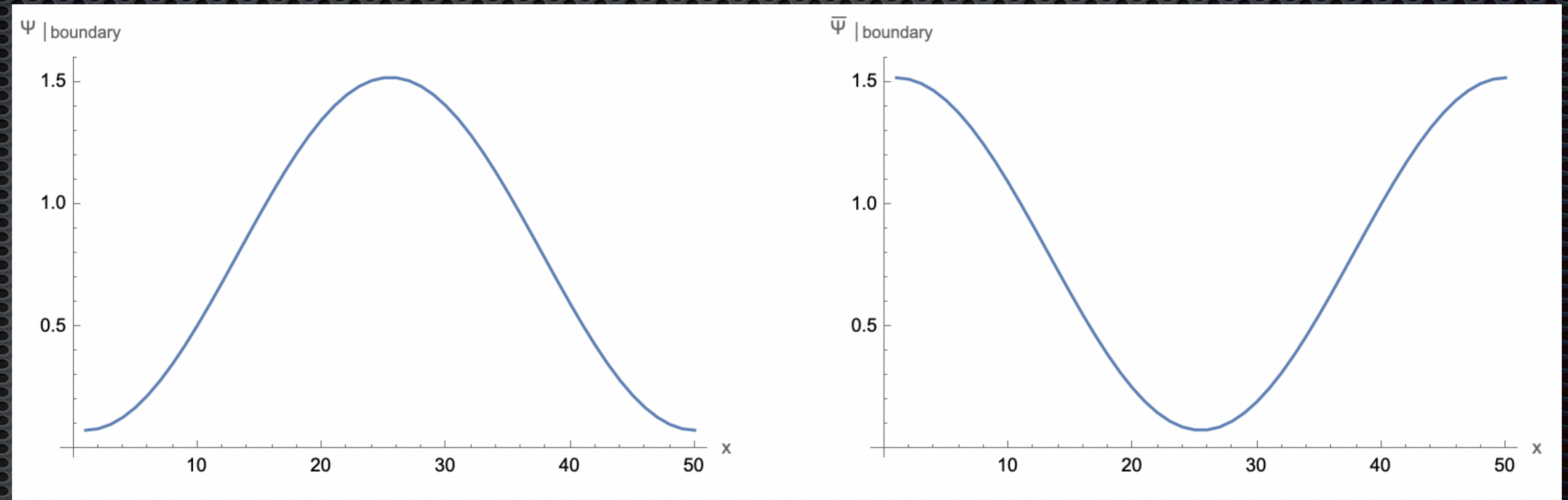
$$\bar{\Psi} \sim s(1 + \xi(x))r$$

# Non-Hermitian inhomogeneous holography

Space dependent non-Hermitian deformation

$$\Psi \sim s(1 - \xi(x))z$$

$$\bar{\Psi} \sim s(1 + \xi(x))z$$

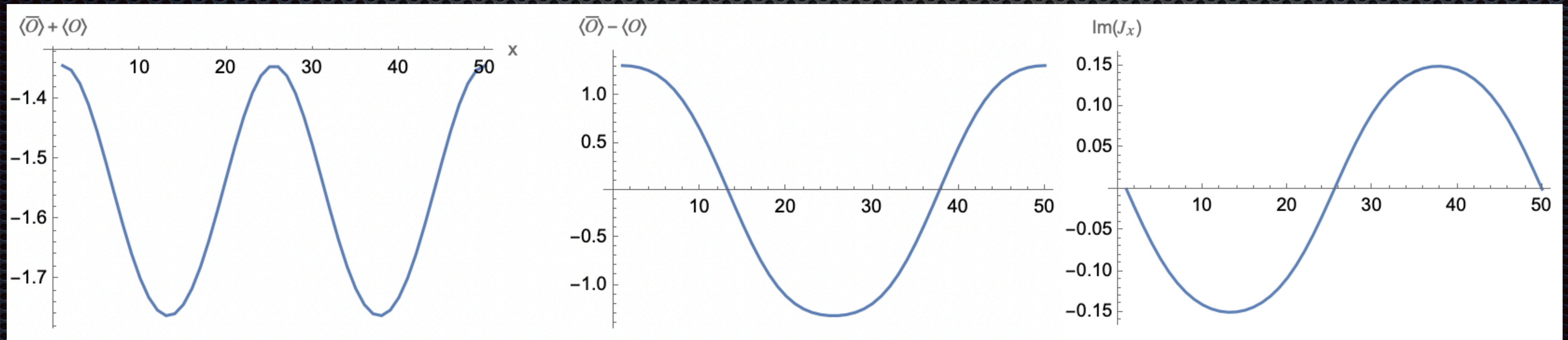


in asymptotically AdS black brane geometry ( $T/s \approx 0.30$ )

[see also Landsteiner&Morales-Tejera'22]

# Non-Hermitian inhomogeneous holography

$$\Psi \sim s(1 - \xi(x)), \quad \bar{\Psi} \sim s(1 + \xi(x)) \quad \implies \quad \langle \text{Im}(J_x) \rangle \neq 0$$



$$T/s \approx 0.30$$

Real geometries dual to non-Hermitian states

# Non-Hermitian inhomogeneous holography

PT-symmetric (pseudo-Hermitian) phase: gauging the Dyson map

[see also Chernodub&Millington'97]

PT-symmetric

$$\Psi \sim s(1 - \xi(x))z$$

$$\bar{\Psi} \sim s(1 + \xi(x))z$$

$\sim$

Hermitian

$$\Psi \sim s\sqrt{1 - \xi(x)^2}z$$

$$\bar{\Psi} \sim s\sqrt{1 - \xi(x)^2}z$$

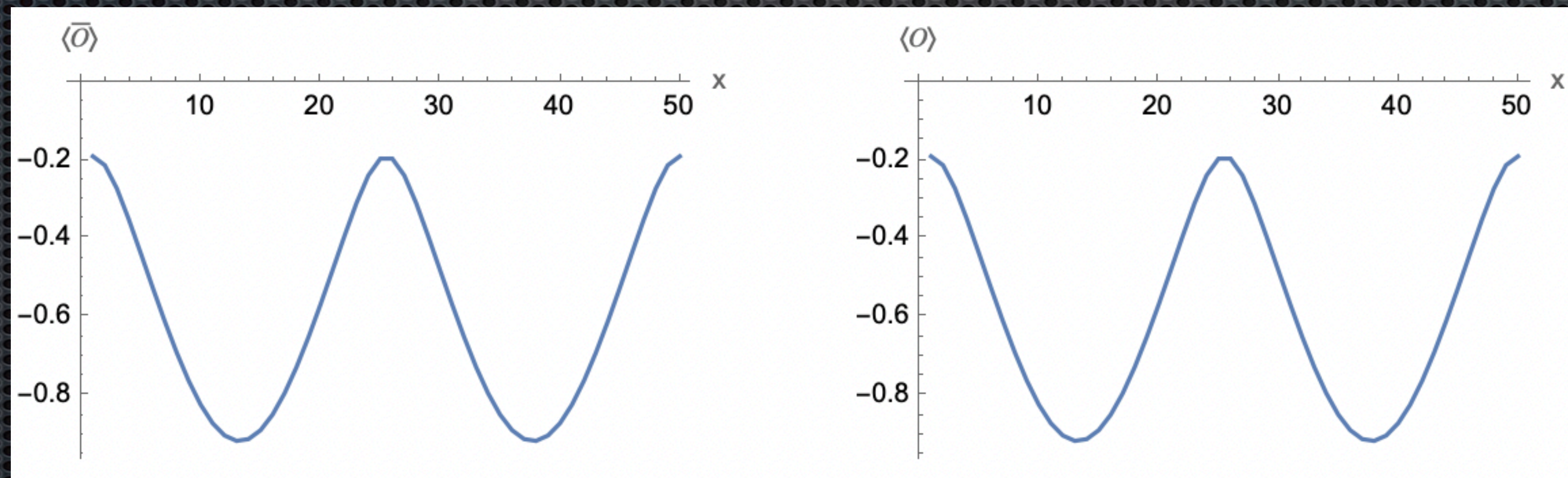
$$A_x = -\frac{i}{2}\partial_x \log\left(\frac{1 - \xi(x)}{1 + \xi(x)}\right)$$



# Non-Hermitian inhomogeneous holography

PT-symmetric (pseudo-Hermitian) phase: gauging the Dyson map

$$\Psi \sim s(1 - \xi(x)), \quad \bar{\Psi} \sim s(1 + \xi(x)), \quad A_x = -\frac{i}{2} \partial_x \log \left( \frac{1 - \xi(x)}{1 + \xi(x)} \right)$$

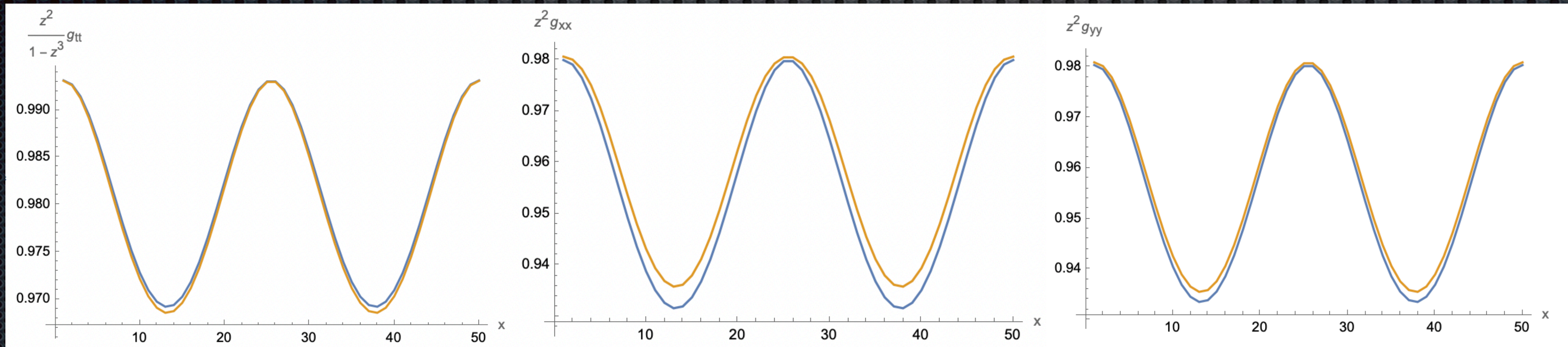


$$\langle J_x \rangle = 0$$

$$T/s \approx 0.30$$

# Non-Hermitian inhomogeneous holography

non-Hermitian  $\langle \text{Im}(J_x) \rangle \neq 0$  vs PT-symmetric geometry



$T/s \approx 0.30$

# Non-hermitian holography

## Overview & To do

- (Minimal) Holographic model of nH PT-symmetric theories
- Exhibits PT-symmetric and PT-broken phases
- Modulated pseudo-Hermitian and non-Hermitian solutions
- QNMs (stability) and transport
- Finite charge density and spontaneous symmetry breaking
- ...

