The quantum entropy of extremal black holes and the Schwarzian theory

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Black holes are quantum-statistical systems

Quantum entropy



Can we calculate $\dim(\mathcal{H}_{\rm BH})$ in gravity?

- Without positing additional discretization of horizon
 [cf Wheeler]
- Microcanonical: Supersymmetric BHs believed to be independent quantum systems



Can we calculate holographic observables at finite N independently in AdS and CFT?



 Supersymmetric indices provide a set of prototype examples: well-defined and calculable on both sides

♦ AdS₂/CFT₁: dim(*H*) of susy BHs (Main example today) [A.Dabholkar, J.Gomes, S.M. '10 -'14; L. Iliesiu, S.M., G.J. Turiaci, '22]

* AdS_{d+1}/CFT_d : similar structure in d = 2, 3, 4

Supersymmetric index in gravity

How does susy BH contribute to the index?

$$W_{\text{grav}}^{\text{index}} = \int Dg_{MN} e^{\mu_i \int \mathcal{A}_i - S_{\text{grav}}(g_{\mu\nu})} Z_{1-\text{loop}}(g_{MN})$$
$${Q\psi_M(g_{MN}) = 0}$$

- Susy BHs are extremal and naively have infinite action.
- Susy BH contribution to index defined as limit $T \rightarrow 0$ of non-extremal susy configurations Cabo-Bizet, Case

(Lo

Ads, $\times S^3$

- complex deformation
- smooth geometry
- smooth limit

Cabo-Bizet, Cassani, Martelli, S.M. '18 in AdS5/CFT4

Cassani, Papini '19
$$\pi i$$
 Bobev, Crichigno '19
in AdS4,6,7, ...

• Same is true for susy BHs in asymptotically flat space [Iliesiu, Kologlu, Turiaci '19] [Boruch, Iliesiu, Murthy, Turiaci, in progress]

 A^R



• We start from AdS_2 limit and then consider T-deformation.

Localization in supergravity

[A.Dabholkar, J.Gomes, S.M. '10 -'14; L. Iliesiu, S.M., G.J. Turiaci, '22]

1. Formalism: Construct rigid supercharge Q in off-shell supergravity (deformation of BRST structure).

[B. de Wit, S.M., V. Reys, '18] [I. Jeon, S.M., '18]

2. All solutions of localization equations $Q \Psi = 0$, w/AdS₂ × S² boundary conditions. ϕ^{I} [R.G

[R.Gupta, S.M. '12]



3. Evaluate full supergravity action on these solutions (include all higher derivative terms). D-terms do not contribute!

[S.M., V.Reys, '13]

4. Compute one-loop determinant.
[Cardoso, de Wit, Mahapatra '12, S.M., V. Reys, '15, Y. Ito, R. Gupta, I. Jeon, '15, I. Jeon, S.M. '18, L.Iliesiu, S.M., G.J. Turiaci, '22]

1/8 BPS BH in N=8 string theory

- Write N=8 string theory as an effective N=2 sugra with 15 vectors, 10 hypers, 6 spin 3/2 multiplets
- Prepotential for vector multiplets is tree-level exact

$$F(X) = -\frac{1}{2} \frac{C_{abc} X^a X^b X^c}{X^0}$$

$$\blacktriangleright W_{\text{grav}}^{(1)}(N) = \int \frac{d\sigma}{\sigma^{9/2}} \exp(\sigma + \pi^2 N/4\sigma) = \widetilde{I}_{7/2}(\pi\sqrt{N})$$

• Microscopic index $W_{micro}(N)$ known from D1-D5 counting. Coefficients of a known modular form.

[Maldacena, Moore, Strominger '00]

Comparison to microscopics

1/8 BPS BH in N=8 string theory

[A.Dabholkar, J.Gomes, S.M. '11]

Ν	$W_{ m micro}(N)$	$W_{\rm grav}^{\rm index(1)}(N)$
3	8	7.97
4	12	12.2
7	39	38.99
8	56	55.72
	152	152.04
12	208	208.45
15	513	512.96
•••	•••	•••
10^5	exp(295.7)	exp(295.7)
$\overline{d_{\text{micro}}(N)} = Z_{\text{grav}}^{\text{index}(1)}(N) \left(1 + O(e^{-\pi\sqrt{N}/2})\right)$		

Can we calculate the exponentially suppressed corrections?



 Treatment of subtle effects coming from taking a BH to zero temperature. Some modes become massless.

[Sachdev '15; Almheiri, Kang '16; Maldacena, Stanford, Yang '16; Moitra, Trivedi, Vishal '18, Maldacena, Turiaci, Yang '19; Ghosh, Maxfield, Turiaci '19; Iliesiu, Turiaci, 20; Iliesiu, Kruthoff, Turiaci, Verlinde, 20; Heydeman, Iliesiu, Turiaci, Zhao, '20;...]

Modular symmetry — Analytic formula (infinite sum) for microscopic degeneracy

[Analytic number theory: G. Hardy, S. Ramanujan ('20); H. Rademacher ('38)]

$$W_{\text{micro}}(N) = \sum_{c=1}^{\infty} c^{-9/2} K_c(N) \widetilde{I}_{7/2} (\pi \sqrt{N}/c)$$

$$= \widetilde{I}_{7/2} (\pi \sqrt{N}) \qquad \text{All order pert.} \\ \text{around } AdS_2$$

$$+ \sum_{c>1}^{\infty} c^{-9/2} K_c(N) \widetilde{I}_{7/2} (\pi \sqrt{N}/c)$$

$$\widetilde{I}_{\rho}(z) = 2\pi (\frac{z}{4\pi})^{-\rho} I_{\rho}(z) \\ = e^z + \dots \quad \text{(I-Bessel function)}$$

$$K_c(N) = \sum_{\substack{0 \le d \le c \\ (c,d)=1}} e^{2\pi i \frac{N}{4} \frac{d+d^{-1}}{c}} M(c,d) \quad \text{(Kloosterman sum)}$$



Zero modes and mass gap

Subtlety in localized path integral

 Zero modes in AdS2 coming from "pure gauge transformations" which don't vanish at boundary. [Camporesi, Higuchi '95]



Euclidean $\mathbf{AdS_2} \times \mathbf{S^2}$

- Effect of zero mode at large charges calculated using ultra-locality [Witten] : important contribution to logarithmic corrections to Bekenstein-Hawking entropy [Sen]
- Regulate zero-modes by turning on a small (susy) temperature.
- Schwarzian mode: spontaneous vs explicit symmetry breaking

Regulate zero modes and then take $T \rightarrow 0$ [L. Iliesiu, S.M., G.J.Turiaci '22]



[..., Stanford, Witten '17; Mertens, Turiaci, Verlinde '17; Heydeman, Iliesiu, Turiaci, Zhao, '20]

Results:

- No mass-gap for generic extremal BHs $Z_{
 m grav}(T,Q) = T^{\alpha} f(Q)$
- For supersymmetric BHs, lpha=0 , mass-gap $1/S_{
 m BH}^{3/2}$
- Volume of space of super-Schwarzian modes on orbifold = 1/c

From index to entropy

Black hole degeneracy = index

- In the microscopic theory, one actually counts the supersymmetric index ${\rm Tr}_{\cal H}\,(-1)^F$. This is protected.
- 4d supersymmetric black holes are spherically symmetric and therefore have zero net angular momentum $J_0 = F$
- AdS_2 geometry \Longrightarrow microcanonical ensemble \Longrightarrow every state has F=0 \Longrightarrow $\operatorname{Tr}_{\mathcal{H}}(-1)^{\mathrm{F}} = \operatorname{Tr}_{\mathcal{H}} 1$

[A. Sen '10; A. Dabholkar, J. Gomes, S.M., A. Sen, '10]

This argument can be extended to Schwarzian modes.
 The result is that only J=0 states contribute.

[Iliesiu, Kologlu, Turiaci '19; L. Iliesiu, S.M., G.J. Turiaci '22]

Higher-dimensional AdS spaces: status

Structure of AdS_{d+1}/CFT_d in d>1

$$\mathcal{I}_N(\tau) \simeq \sum_{\substack{d \\ c \in \mathbb{Q}}} \exp\left(\frac{1}{c} F_{\rm BH}^{\rm grav}(\tau - d/c)\right)$$

- Family of saddles CFT4 [Benini, Milan'18] [Cabo-Bizet, S.M.'19] [...] of CFT index. CFT3 [Benetti-Genolini, Cabo-Bizet, S.M.,'23]
- Family of gravitational AdS3 [Maldacena, Strominger.,'98]
 saddles (orbifolds in string/M-theory)
 AdS4 [Benetti-Genolini, Cabo-Bizet, S.M.,'23]
 AdS5 [Aharony, Benini, Mamroud, Milan,'21]
- 1/N Perturbation expansion around each saddle. Hints.
 AdS5 [Aharony, Benini, Mamroud, Milan,'21] AdS3 [Ciceri, Jeon, S.M.,'23]
 AdS4 [Pando Zayas, Xin,'20] [Mamroud '22]
 Bobev, Reys, + Charles, Hristov, Choi '21, '22]
- Number theory CFT2 [Dijgraaf-Maldacena-Moore-Verlinde '00][Maschot-Moore '10] structure. CFT4 [Cabo-Bizet, S.M.'19] [Garoufalidis, S.M., Zagier, ??]

Thank you for your attention!