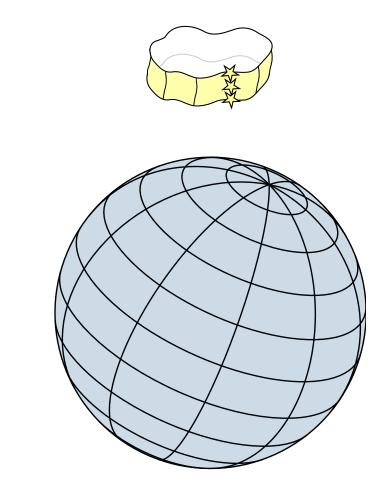
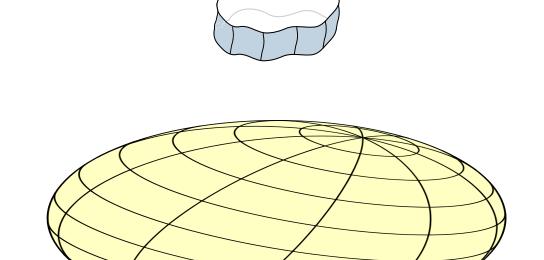
Eurostrings 2023, Gijón/Xixón 28th of April

Jordanian deformations of the

AdS₅ × S⁵ superstring





Sibylle Driezen

based on 2112.12025; 2207.14748; 2212.11269 with

Riccardo Borsato; J. Luis Miramontes; Juan Miguel Nieto García and Leander Wyss



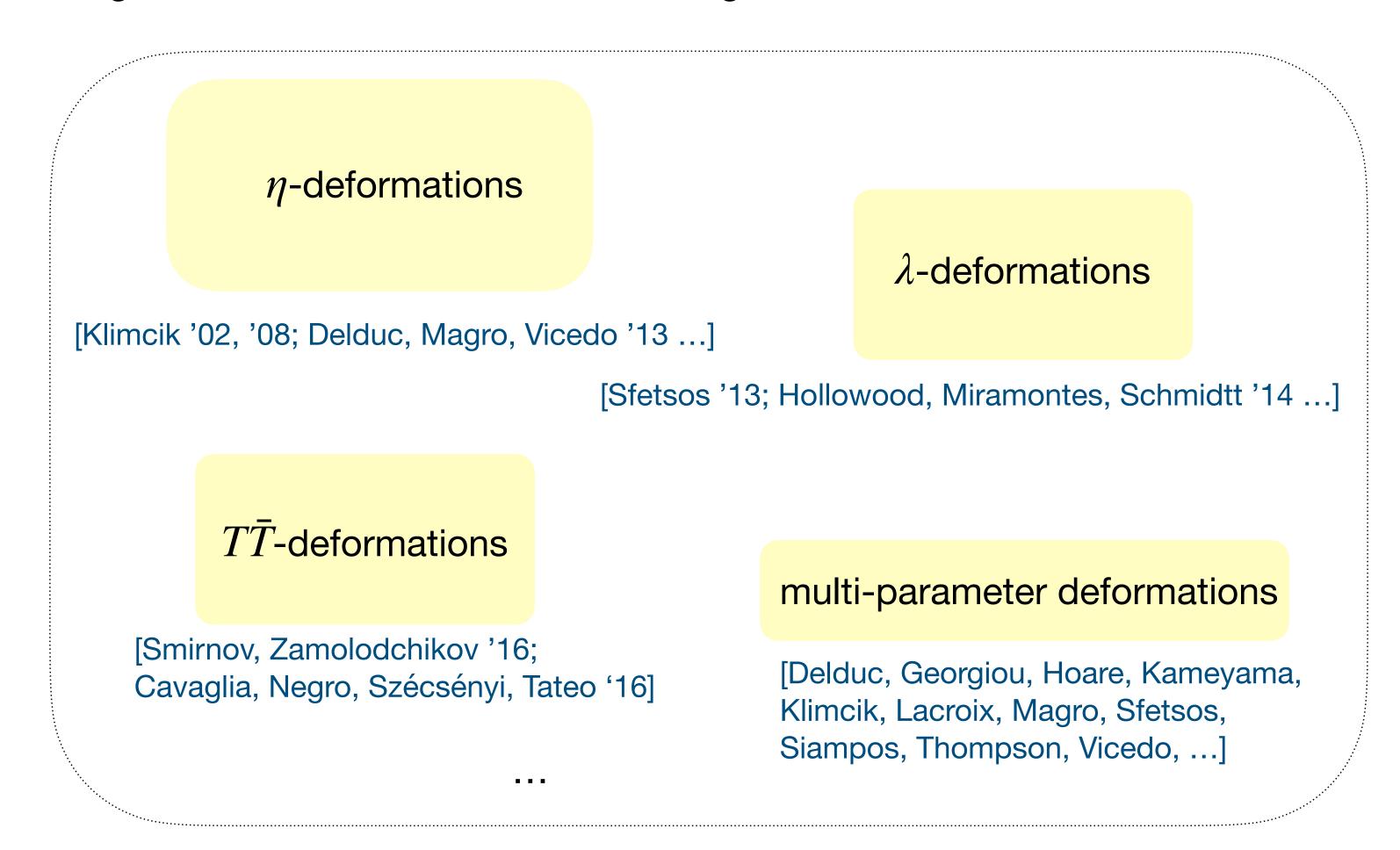
Why integrable deformations of worldsheet sigma-models

Integrability program proven to be very successful for strongly coupled gauge-theories and strings in background fields

- → tailored toolbox of analytic techniques (based on hidden symmetries) to derive exact results
- \rightarrow rigid proofs for e.g. $AdS_5 \times S^5 \leftrightarrow \mathcal{N} = 4$ SYM

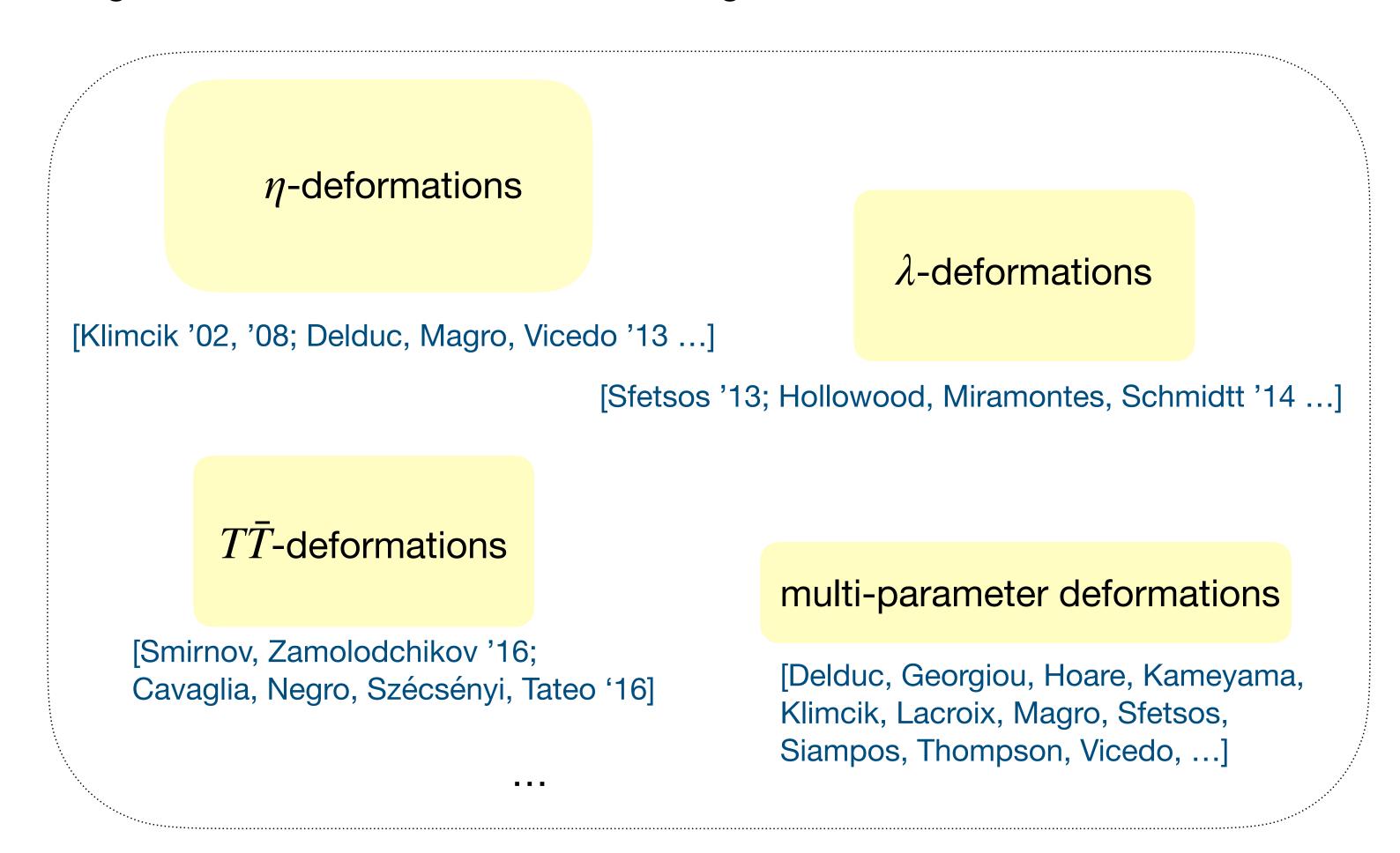
- ◆ holography: deform to AdS/CFT systems with less manifest symmetries (non-conformal, non-maximally susy)
 whilst preserving the ability to apply integrability program
- → more formally: generalise the integrable toolbox (underlying hidden structures at play)
 - whilst being rare: explore family of integrable models and their mathematical formulation

Today: whole zoo of integrable deformations of worldsheet sigma-models



a lot known about their SUGRA embedding [Borsato, Demulder, SD, Hassler, Hoare, Seibold, Sfetsos, Tseytlin, Thompson, Van Tongeren, Wulff,...]

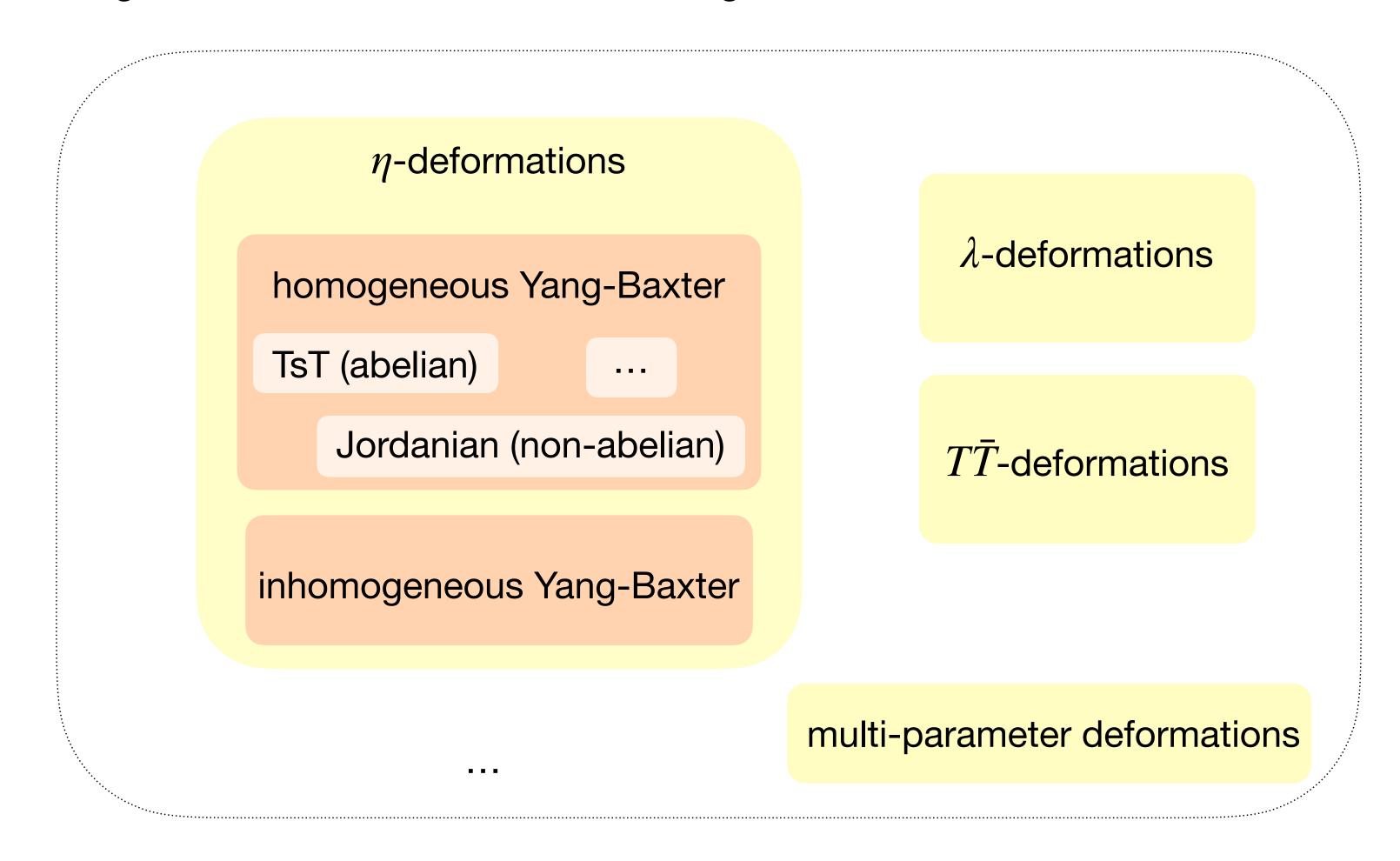
Today: whole zoo of integrable deformations of worldsheet sigma-models



a lot known about their SUGRA embedding

[Borsato, Chervonyi, Demulder, SD, Hassler, Hoare, Lunin, Seibold, Sfetsos, Tseytlin, Thompson, Van Tongeren, Wulff,...]

Today: whole zoo of integrable deformations of worldsheet sigma-models



a lot known about their SUGRA embedding

Extension of integrability methods developed for $AdS_5 \times S^5 \leftrightarrow \mathcal{N} = 4$ SYM to deformations?

spectral problem known from first-principles for restricted class of "diagonal" TsT

[Beisert, Roiban '05; Gromov, Levkovich-Maslyuk '10; de Leeuw, Van Tongeren '12; Kazakov '18...]

based on reformulation of deformed model as undeformed model with local twisted boundary conditions

[Frolov, Roiban, Tseytlin '05; Alday, Arutyunov, Frolov '05]

 η -deformations

homogeneous Yang-Baxter

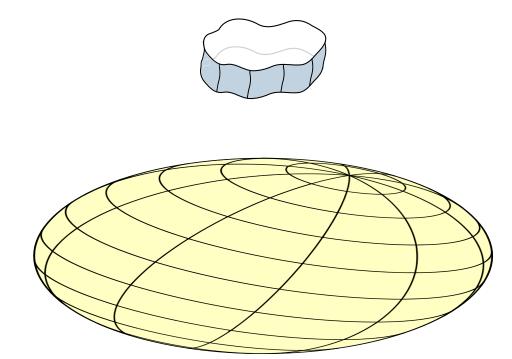
TsT (abelian)

Jordanian (non-abelian)

inhomogeneous Yang-Baxter

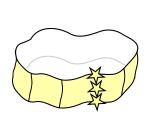
Outline

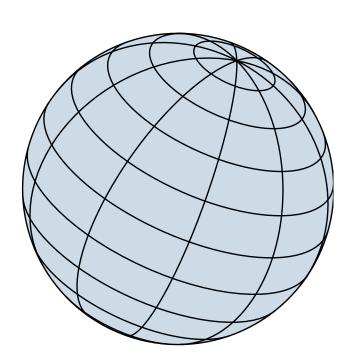
◆ Brief intro to Homogeneous Yang-Baxter deformations



◆ Reformulation in terms of undeformed twisted models

- → Jordanian deformations and why we care
- \bullet Classification of Jordanian deformations for ${\rm AdS}_5\times {\rm S}^5$ preserving preserving D=10 IIB SUGRA
- ◆ Spectral curve methods





Homogeneous Yang-Baxter (HYB) deformations

$$g = g(X^{\mu}(\sigma^{\alpha})) \in G = \text{Lie}(\mathfrak{g})$$

$$S = \int d^2\sigma \langle \partial_+ gg^{-1}, (1 - \eta R)^{-1} \partial_- gg^{-1} \rangle$$

HYB of Principal Chiral Model

 $\eta \in \mathbb{R}$

$$R: \mathfrak{g} \to \mathfrak{g}: T_a \to R(T_a) = R_a^b T_b$$

$$[Rx, Ry] - R([Rx, y] + [x, Ry]) = 0$$

$$R^T = -R$$

$$R^{ab}f_{ab}^{c} = 0$$

classical Yang-Baxter equation

antisymmetric

unimodularity

- Preserves worldsheet integrability [Klimcik '02, '08; Delduc, Magro, Vicedo '13] and SUGRA [Borsato, Wulff '16]
- TsT = "abelian-R" (Im(R) is abelian subalgebra)
- $-\operatorname{Im}(R)$ is **non-abelian** subalgebra of \mathfrak{g}

PCM

 $(\mathsf{CYBE} + R^T = -R)$

Currents

Conserved (EOM) & Flat

Flat Lax connection

 $\tilde{J}_{\pm} = \tilde{g}^{-1} \partial_{\pm} \tilde{g}$ $d \star \tilde{J} \approx 0 \quad \& \quad d\tilde{J} + \tilde{J} \wedge \tilde{J} = 0$

$$\mathscr{L}_{\pm}(z) = \frac{\tilde{J}_{\pm}}{1 \pm z}, \quad z \in \mathbb{C}$$

Currents

Conserved (EOM) & Flat

Flat Lax connection

$$A_{\pm} = \text{Ad}_g^{-1} (1 \pm \eta R)^{-1} \partial_{\pm} g g^{-1}$$

HYB

$$d \star A \approx 0 \& dA + A \wedge A \approx 0$$

$$\mathscr{L}_{\pm}(z) = \frac{A_{\pm}}{1 \pm z}, \quad z \in \mathbb{C}$$

(hallmark of classical integrability)

PCM

HYB

$$(\mathsf{CYBE} + R^T = -R)$$

Currents

 $\tilde{J}_{\pm} = \tilde{g}^{-1} \partial_{\pm} \tilde{g}$ $d \star \tilde{J} \approx 0 \quad \& \quad d\tilde{J} + \tilde{J} \wedge \tilde{J} = 0$ Conserved (EOM) & Flat

Conserved (EOM) & Flat $d \star A \approx 0$ & $dA + A \wedge A = 0$

 $A_{\pm} = \mathrm{Ad}_{g}^{-1} (1 \pm \eta R)^{-1} \partial_{\pm} g g^{-1}$

Flat Lax connection

$$\mathscr{L}_{\pm}(z) = \frac{\tilde{J}_{\pm}}{1 \pm z}, \quad z \in \mathbb{C}$$

Flat Lax connection

Currents

$$\mathscr{L}_{\pm}(z) = \frac{A_{\pm}}{1 \pm z}, \quad z \in \mathbb{C}$$

Exploit on-shell equivalence:

$$\tilde{J}_{\pm}[\tilde{g}] \approx A_{\pm}[g]$$

$$\tilde{J}_{\pm}[\tilde{g}] \approx A_{\pm}[g]$$
 \Rightarrow relate $g = \mathcal{F}\tilde{g}$

PCM

HYB

$$(\mathsf{CYBE} + R^T = -R)$$

Currents

 $d \star \tilde{J} \approx 0 \quad \& \quad d\tilde{J} + \tilde{J} \wedge \tilde{J} = 0$ Conserved (EOM) & Flat

Flat Lax connection

$$\mathscr{L}_{\pm}(z) = \frac{\tilde{J}_{\pm}}{1+z}, \quad z \in \mathbb{C}$$

 $\tilde{J}_{\pm} = \tilde{g}^{-1} \partial_{\pm} \tilde{g}$

Currents

Flat Lax connection

$$A_{\pm} = \mathrm{Ad}_{g}^{-1} (1 \pm \eta R)^{-1} \partial_{\pm} g g^{-1}$$

Conserved (EOM) & Flat $d \star A \approx 0$ & $dA + A \wedge A = 0$

$$\mathscr{L}_{\pm}(z) = \frac{A_{\pm}}{1 \pm z}, \quad z \in \mathbb{C}$$

Exploit on-shell equivalence:

$$\tilde{J}_{\pm}[\tilde{g}] \approx A_{\pm}[g]$$

$$\tilde{J}_{\pm}[\tilde{g}] \approx A_{\pm}[g]$$
 \Rightarrow relate $g = \mathcal{F}\tilde{g}$

deformed periodic HYB

$$g(2\pi) = g(0)$$

PCM

HYB

$$(\mathsf{CYBE} + R^T = -R)$$

Currents

 $\tilde{J}_{+} = \tilde{g}^{-1} \partial_{\pm} \tilde{g}$

 $A_{\pm} = \mathrm{Ad}_{g}^{-1} (1 \pm \eta R)^{-1} \partial_{\pm} g g^{-1}$

Conserved (EOM) & Flat

 $d \star \tilde{J} \approx 0 \quad \& \quad d\tilde{J} + \tilde{J} \wedge \tilde{J} = 0$

 $d \star A \approx 0 \& dA + A \wedge A = 0$

Flat Lax connection

$$\mathscr{L}_{\pm}(z) = \frac{\tilde{J}_{\pm}}{1 \pm z}, \quad z \in \mathbb{C}$$

Flat Lax connection

Conserved (EOM) & Flat

Currents

$$\mathscr{L}_{\pm}(z) = \frac{A_{\pm}}{1 \pm z}, \quad z \in \mathbb{C}$$

Exploit on-shell equivalence:

$$\tilde{J}_{\pm}[\tilde{g}] \approx A_{\pm}[g]$$

$$\tilde{J}_{\pm}[\tilde{g}] \approx A_{\pm}[g]$$
 \Rightarrow relate $g = \mathcal{F}\tilde{g}$

undeformed twisted PCM

 $\tilde{g}(2\pi) = W\tilde{g}(0) = \mathcal{F}^{-1}(2\pi)\mathcal{F}(0)\tilde{g}(0)$

deformed periodic HYB

$$g(2\pi) = g(0)$$

non-local expression W[g]

[Matsumoto, Yoshida '15; Vicedo '15; Van Tongeren '18]

: local and closed expression [Borsato, SD, Miramontes '21] $W[\tilde{g}]$

HYB-picture

PCM-picture

$$g(2\pi) = g(0) \qquad \Rightarrow \qquad \tilde{g}(2\pi) = W\tilde{g}(0),$$

$$W = e^{\eta RQ[\tilde{g}]}$$

$$\partial_{\tau}W = 0$$
on-shell
$$\approx$$
equivalent

Jordanian deformations

When $\mathfrak{Sl}(2,R)$ subalgebra of isometries

$$[h, e_{\pm}] = \pm e_{\pm}$$
, $[e_{+}, e_{-}] = 2h$ \Rightarrow $R(h) = \frac{e_{+}}{2}$, $R(e_{-}) = -h$, $R(e_{+}) = 0$

- truly **non-abelian**: $Im(R) = \{h, e_+, + \text{ possibly supercharges}\} \Rightarrow \text{ goes beyond TsT}$
- have a twist that is always diagonal $[\leftrightarrow]$ abelian (TsT) & almost-abelian (sequence of TsT's)]

[Borsato, SD, Miramontes '21]

→ important for **usability of twisted models** in semi-classical integrable methods

Why

$$\Omega(z) = W^{-1} \operatorname{Pexp} \left(-\int \mathcal{L}(z) \right)$$

$$= W^{-1} \left(1 + zQ \right) + \mathcal{O}(z^2)$$
has
$$\partial_{\tau} \lambda(z) = 0$$

- ♦ has local asymptotics ⇒ local charges after going to diagonal or Jordan form
- lacktriangle has diagonalisable asymptotics \Rightarrow allows to reconstruct eigenvalues of $\Omega(z)$ on $\mathbb C$ for any sol.

Long been "ignored" because **bosonic** R is not unimodular, i.e. $R^{ab}f_{ab}^{\ \ c} \neq 0 \Rightarrow \text{gSUGRA}$,

until [Van Tongeren '19] showing inclusion of supercharges in ${\rm Im}(R)$ can mitigate this, i.e. $\tilde{R}^{ab}f_{ab}^{\ \ c}=0 \Rightarrow {\rm SUGRA}$

Classifying Jordanian models of $AdS_5 \times S^5$

All (unimodular) Jordanian deformations of $AdS_5 \times S^5$ superstring, $\mathfrak{g} = \mathfrak{psu}(2,2 \mid 4)$, classified in [Borsato, SD '22]

$$R = R^{ab}T_a \wedge T_b = h \wedge e_+ - \sum_{i=1}^N e_{-i} \wedge e_i \qquad \text{w.} \qquad \begin{array}{c} [h, e_+] = e_+, & [h, e_{\pm i}] = (1/2 \pm \xi_i) \ e_{\pm i}, \\ [e, e_{\pm i}] = 0, & [e_i, e_j] = \delta_{i,-j}e \end{array}$$

$$\text{CYBE (integrability) [Tolstoy '04]}$$

$$\Rightarrow R^{ab}[T_a, T_b] = 4(1 + N_0 - N_1) \ e_+ \qquad \Rightarrow \qquad \text{unimodularity (SUGRA) if } N_1 = N_0 + 1$$

$$\text{\# bosonic extensions} \qquad \text{\# fermionic extensions}$$

Simplest case:

$$N_0 = 0,$$
$$N_1 = 1$$

Isometries max. bosonic supercharges $g \rightarrow g_L g \qquad \begin{array}{c} 5+9 \\ 3+4 \\ 3+9 \\ 6 \\ 5+4 \\ 5+9 \end{array} \qquad \begin{array}{c} 0 \\ 4 \\ 6 \\ 5+4 \\ 5+9 \end{array}$

Classifying Jordanian models of $AdS_5 \times S^5$

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Simplest case:

$$N_0 = 0,$$

$$N_1 = 1$$

Isometries	max. bosonic	supercharges
$g o g_L g$ $R = \mathrm{Ad}_{g_L}^{-1} R \mathrm{Ad}_{g_L}, \mathrm{ad}_{T_{\bar{a}}} R = R \mathrm{ad}_{T_{\bar{a}}}$	5 + 9 3 + 4 3 + 9 5 + 4 5 + 9	0 4 6 8

Spectral Curve methods

[Borsato, SD, Nieto, Wyss '22]

$$\Omega(z) = W^{-1} \operatorname{Pexp}\left(-\int \mathcal{L}(z)\right) = W^{-1}(1+zQ) + \mathcal{O}(z^2)$$

has conserved eigenvalues $\lambda(z)=e^{ip(z)}$, quasimomenta $p(z)=\mathrm{diag}\{\hat{p}_1(z),\hat{p}_2(z),\hat{p}_3(z),\hat{p}_4(z)\,|\,|\,\tilde{p}_1(z),\tilde{p}_2(z),\tilde{p}_3(z),\tilde{p}_4(z)\}$

AdS₅
$$\sim \frac{SU(2,2)}{Sp(1,1)}$$
 S⁵ $\sim \frac{SU(4)}{Sp(2)}$

$$S^5 \sim \frac{SU(4)}{Sp(2)}$$

Expansion around $z \sim 0$ to $\mathcal{O}(z)$ gives **local conserved charges**

$$\hat{p}_1(z) \sim -(E + Q_{\Theta})z + \mathcal{O}(z^2)$$

$$\hat{p}_2(z) \sim -\frac{i}{2} \mathbf{Q}_W + Q_{\Theta} z + \mathcal{O}(z^2)$$

$$\hat{p}_4(z) \sim (E - Q_{\Theta})z + \mathcal{O}(z^2)$$

$$\hat{p}_3(z) \sim +\frac{i}{2}\mathbf{Q}_W + Q_{\Theta}z + \mathcal{O}(z)$$

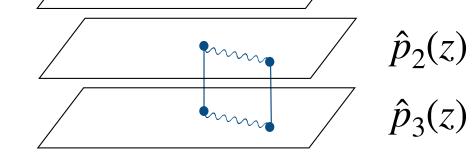
Evaluate $\Omega(z)$ and its eigenvalues $\lambda(z)=e^{ip(z)}$ on simplest non-trivial classical solution for field g

→ point-like in deformed picture "BMN-like"

$$T=\tau,$$
 ...

$$\hat{p}_1(z) = -\hat{p}_4(z) = \frac{2\pi\sqrt{z^2 - \eta^2}}{z^2 - 1},$$

$$\hat{p}_2(z) = -\hat{p}_3(z) = \frac{2\pi z\sqrt{1 - z^2\eta^2}}{z^2 - 1}$$



hum

 $\hat{p}_1(z)$

 $\hat{p}_4(z)$

$$E = 1, \quad \mathbf{Q}_W = 4\pi\eta, \quad Q_{\Theta} = 0$$

Semi-classical corrections to the curve

Method of [Gromov, Vieira '07]: introduce quantum excitations in the form of microscopic cuts ~ quantum poles

$$p_i \rightarrow p_i + \delta p_i$$

Must fulfil number of analytical properties from the curve and from PSU(2,2|4)

- gluing conditions on cuts
- behaviour around poles of the Lax
- \mathbb{Z}_4 symmetry of PSU(2,2|4)
- asymptotic behaviour around $z \sim 0$

→ restrictive enough to fix corrections completely

$$E_{1-\mathrm{loop}} = E_{1-\mathrm{loop}}\left(\eta\right)$$
 (6.39) and $\delta\mathbf{Q}_W = 0$

[Borsato, SD, Nieto, Wyss '22]

analogous to diagonal TsT / β -twisted SYM [Beisert, Roiban '05; de Leeuw, Van Tongeren '12...]

Conclusions and outlook

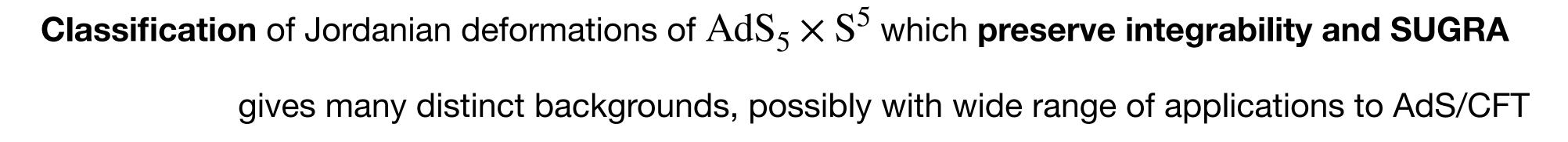
Homogeneous Yang-Baxter deformations as undeformed yet twisted models

allows the use of the **classical spectral curve** method and its semi-classical quantisation unambiguously when the twist is **diagonal**



applied to a particular Jordanian deformation of $AdS_5 \times S^5$ to extract **1-loop** correction to energy of a BMN-like solution

relation and matching to (possible twisted spin-chain of) deformed N=4 SYM side? start up quantum integrable program and match with exact worldsheet results?



find well-behaved solutions and use integrability to study non-symmetric backgrounds?

