

Holographic Torus Entanglement and its RG flow

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William Witczak-Krempa



Based on:

PB, Witczak-Krempa, [arXiv:1611.01846](#)

Outline

1- Introduction

1.1- EE in (3d) QFTs

1.2- Holographic EE

2- Torus Entanglement

2.1- General results & holography

2.2- Renormalized EE

2.3- Holographic RG flow

3- Summary

1- Introduction

Entanglement Entropy

Definition

Starting point: bipartite quantum system $\longleftrightarrow \mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$

Choose some state $\rho \in \mathcal{H}$

Integrate out d.o.f. in the complement $\longrightarrow \rho_A = \text{Tr}_B \rho$

Compute Von Neumann entropy of ρ_A

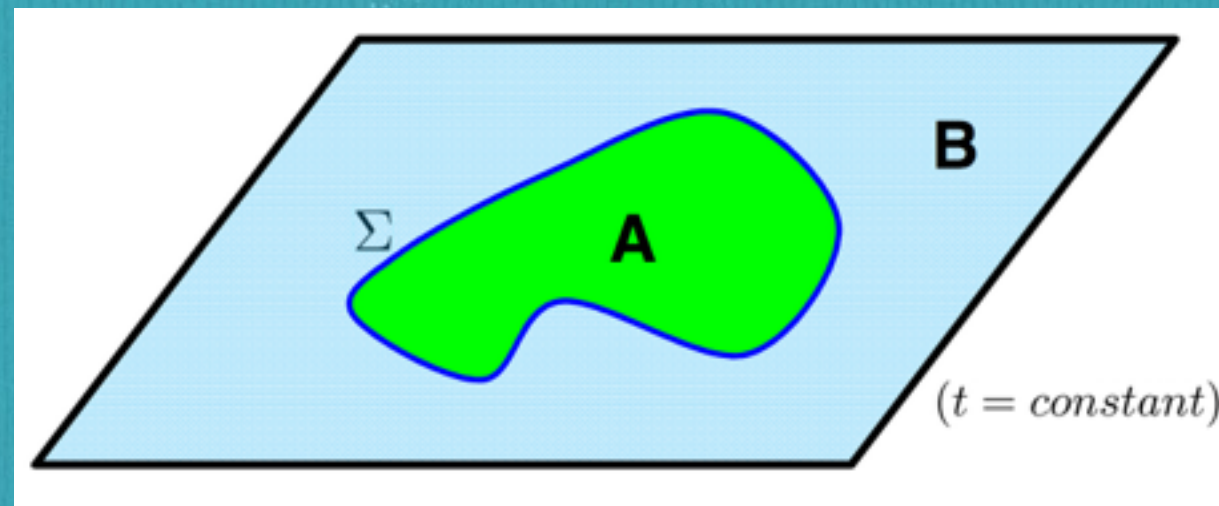
Entanglement Entropy

$$S_{EE} = -\text{Tr} [\rho_A \log \rho_A]$$

Counts # of entangled bits between A and B

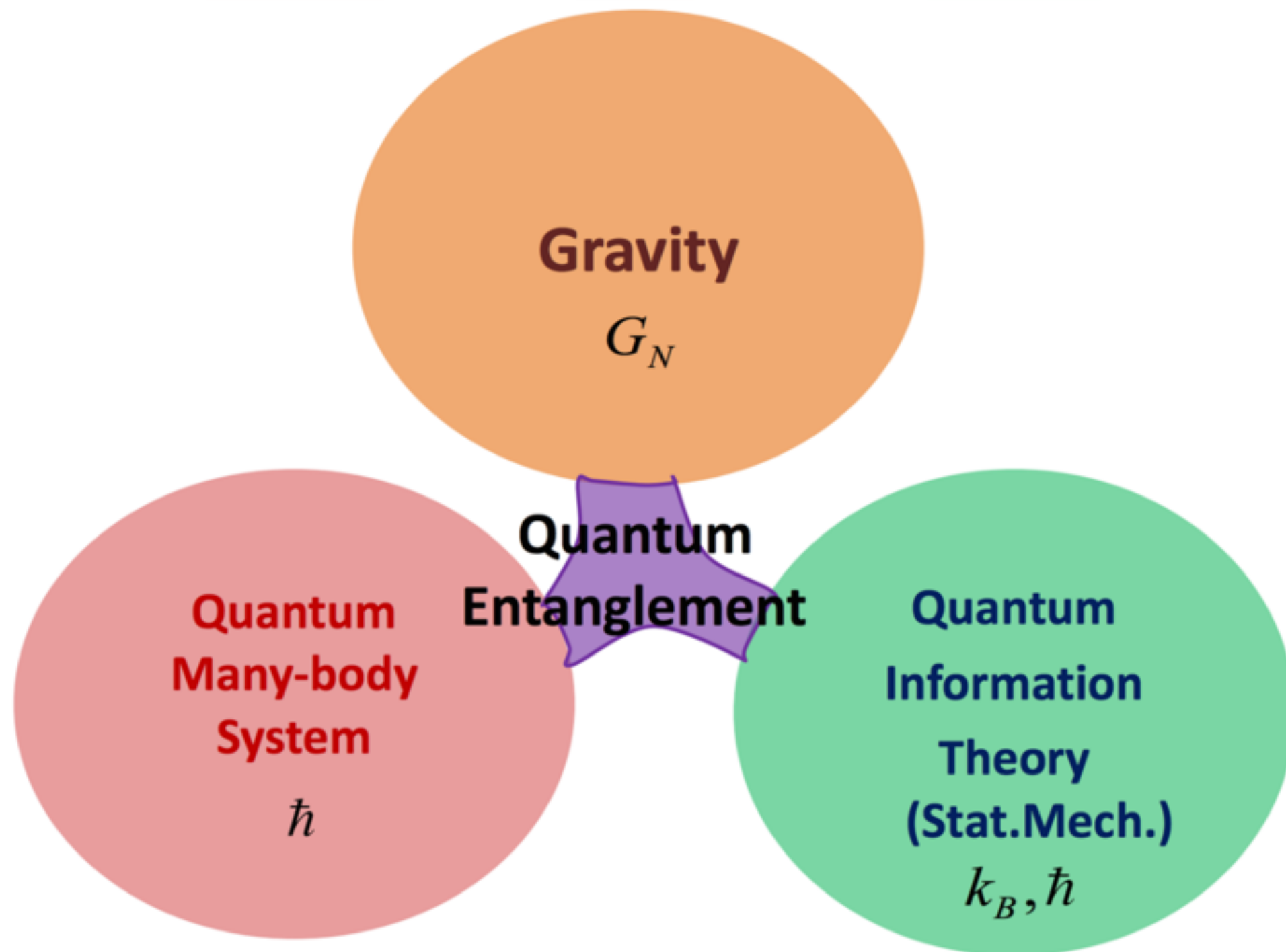
Geometric entropy $\longleftrightarrow \mathbf{A}$ is a spatial region

Focus in this talk



Entanglement Entropy

Why?



1.1- Entanglement Entropy in (3d) QFTs

Entanglement Entropy

In QFTs

$\text{QFT}_d \longrightarrow$ UV divergent but **universal terms**

$$S_{EE}(A) = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_1 \frac{R^{d-3}}{\delta^{d-3}} + \dots$$

regulator
dependent!

Leading term controlled by local correlations at both sides of Σ : **Area law**

(Volume law for
excited states)

Some of the subleading terms encode well-defined information about the theory

Entanglement Entropy

In 3d CFTs

General structure:

“Area”-law term

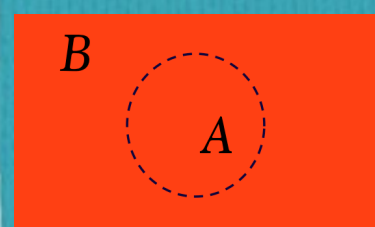
Non-universal constant

Relevant IR scale

Universal term

UV cut-off

$$S_{3d} = \mathcal{B} \frac{L}{\epsilon} + S_{\text{univ}}$$



Disk region

$$S_{\text{univ}} = -F$$

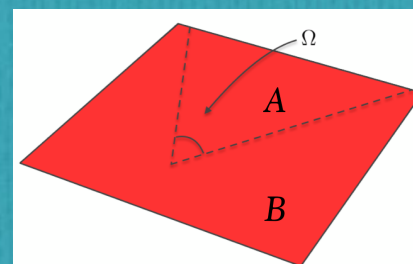
pure constant

equals Free energy on S^3
Casini, Huerta, Myers

F-theorem!

Casini, Huerta

3d version of “c-theorem”



Corner region

$$S_{\text{univ}} = -q(\Omega) \log(L/\epsilon)$$

controlled by local correlator

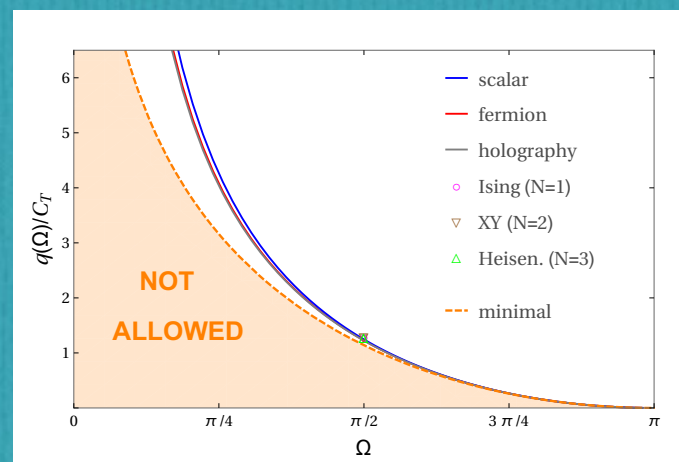
$$\langle T_{ab}(x) T_{cd}(0) \rangle = C_T \frac{\mathcal{I}_{ab,cd}}{x^6}$$

exact in almost smooth limit

$$q(\Omega \sim \pi) = \sigma \cdot (\pi - \Omega)^2$$

$$\frac{\sigma}{C_T} = \frac{\pi^2}{24}$$

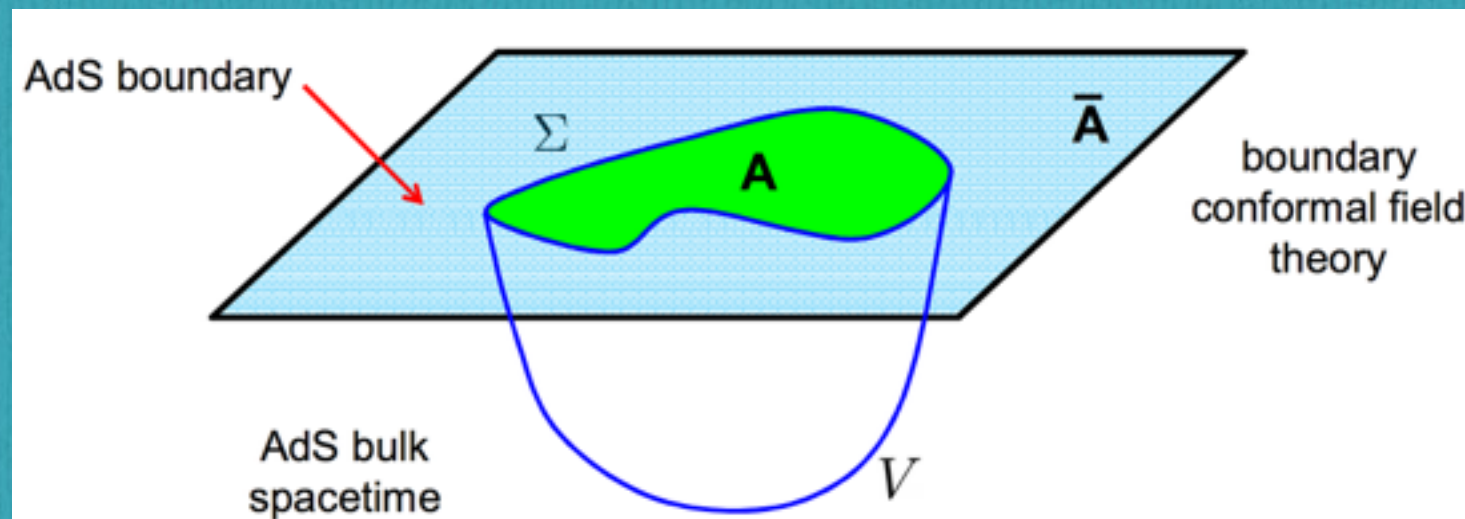
PB, Myers, Witczak-Krempa



1.2- Holographic Entanglement Entropy

Holographic Entanglement Entropy

Ryu-Takayanagi prescription
for CFTs dual to **Einstein gravity**



Ryu Takayanagi +1000 citations!

$$S_{EE}(A) = \min_{V \sim A} \left[\frac{\mathcal{A}(V)}{4G} \right]$$

Extremize area functional $\mathcal{A}(V)$ over all bulk surfaces V whose boundary coincides with Σ

Evaluate $\mathcal{A}(V)$ on the extremal V

Tons of consistency checks and applications *Many people*

Proof

Generalized to: time dependence *Hubeny, Rangamani, Takayanagi*

Lewkowycz, Maldacena

higher-order gravities *Many people*

quantum corrections, etc. *Faulkner, Lewkowycz, Maldacena*

2- Torus Entanglement

Torus Entanglement

Smooth curved surfaces are not ideal for numerics

Pixelization leads to corners \longrightarrow pollutes the result with log. terms

Finite-size effects also pollute unless total space much larger than A

Alternative



Flat but finite entangling surfaces

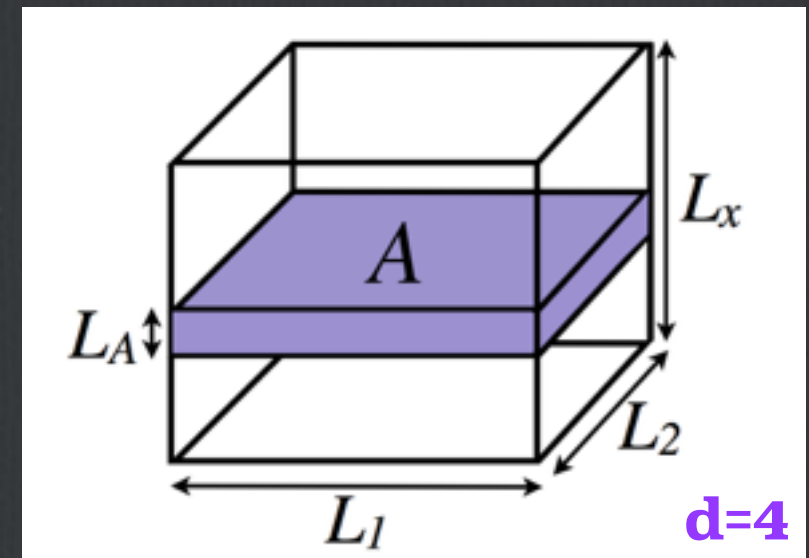
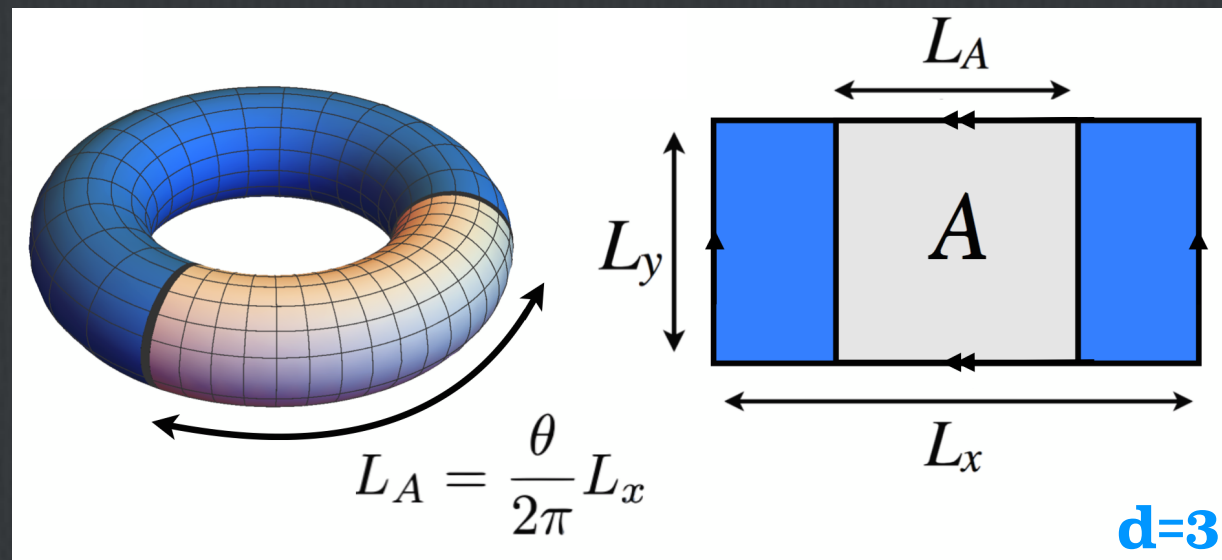


Spaces with non-trivial topology

less explored

2.1- General results & holography

Torus Entanglement



Spatial dimensions form a $T^{(d-1)}$

$$S(A) = B \frac{\text{Area}(\partial A)}{\epsilon^{d-2}} - \chi(\theta; b_i) + \dots \quad b_i = L_x / L_i$$

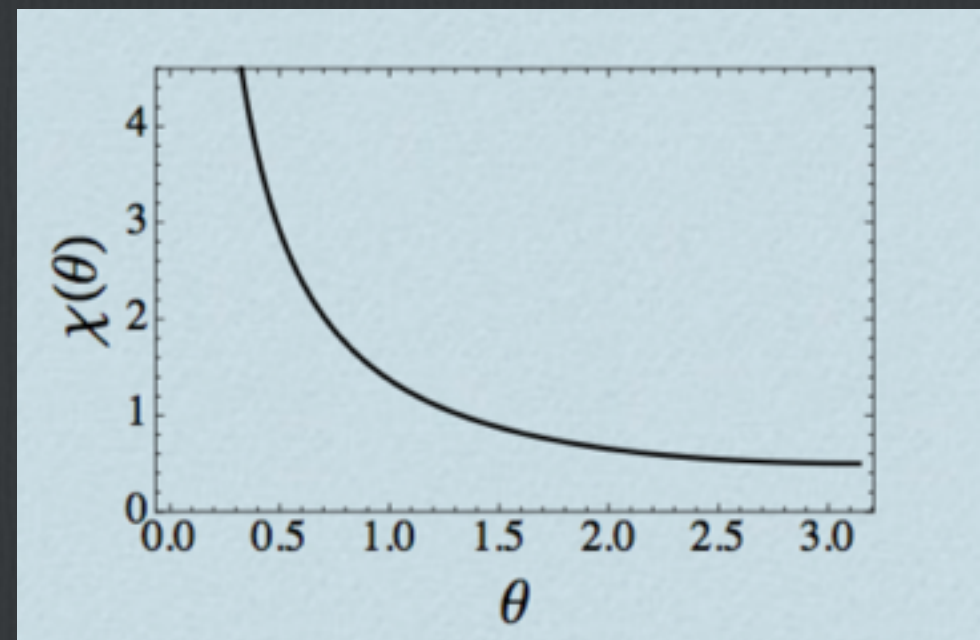
Regulator independent

SSA

$$\chi'(\theta) \leq 0, \quad \chi''(\theta) \geq 0,$$

$$\chi(\theta \rightarrow 0) = \frac{(2\pi)^{d-2} \kappa}{\theta^{d-2} b_1 \dots b_{d-2}}$$

$$\chi(\theta \approx \pi) = \sum_{\ell=0} c_{\ell} \cdot (\pi - \theta)^{2\ell}$$



Witczak-Krempa, Hayward, Melko

Chen, Cho, Faulkner, Fradkin

Torus Entanglement

Holography $\xrightarrow{(T=0)}$ AdS solitons

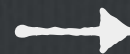
Witten;
Horowitz, Myers

$$ds^2 = \frac{1}{z^2} \left[\frac{dz^2}{f(z)} + f(z) dx^2 + d\vec{y}_{(d-2)}^2 - dt^2 \right]$$

$f(z) = 1 - (z/z_h)^d$ doubly-Wick-rotated black branes

$L_x = 4\pi z_h/d$ smallest dim. must be compact

If additional (d-2) spatial
dimensions also periodic



**conformal boundary
foliated by $\mathbf{T}^{(d-1)}$**

AdS solitons dominate the semiclassical partition function at small temperatures

Focus in this
talk: $d=3$

Torus Entanglement

Chen, Cho, Faulkner, Fradkin
PB, Witczak-Krempa

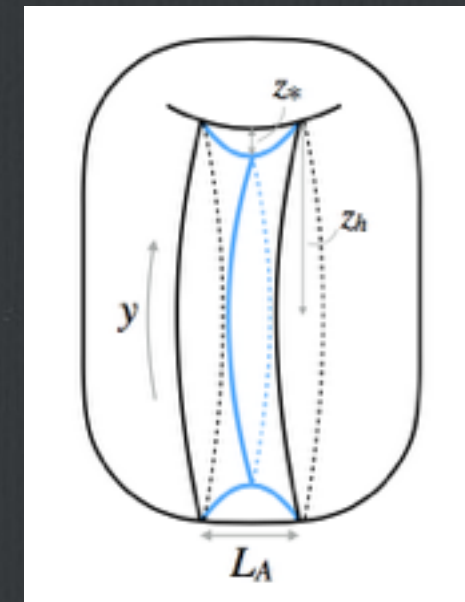
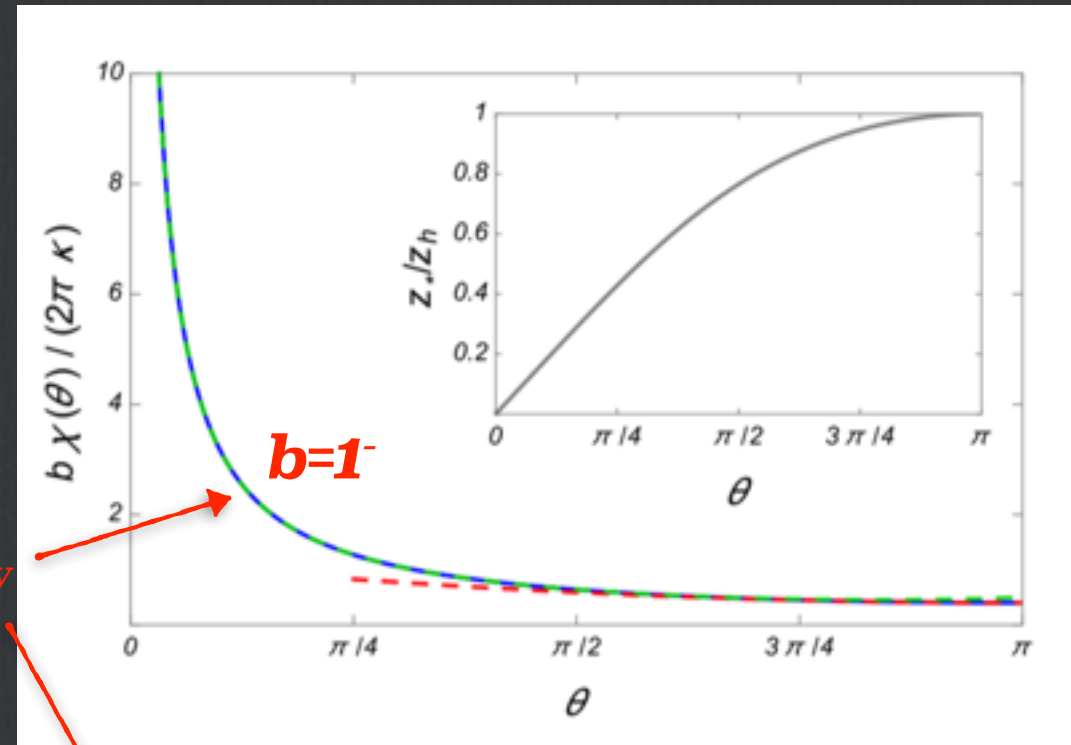
$$L_y > L_x$$

$$\chi(\theta) = \left[\frac{2\pi\kappa}{b} \right] \frac{1}{\theta} + \left[\frac{\Gamma(\frac{1}{4})^{12}}{1306368\pi^7 b G} \right] \theta^5 + \dots$$

Corner like

(Non-generic) jump at $L_x=L_y$

Chen, Cho, Faulkner, Fradkin



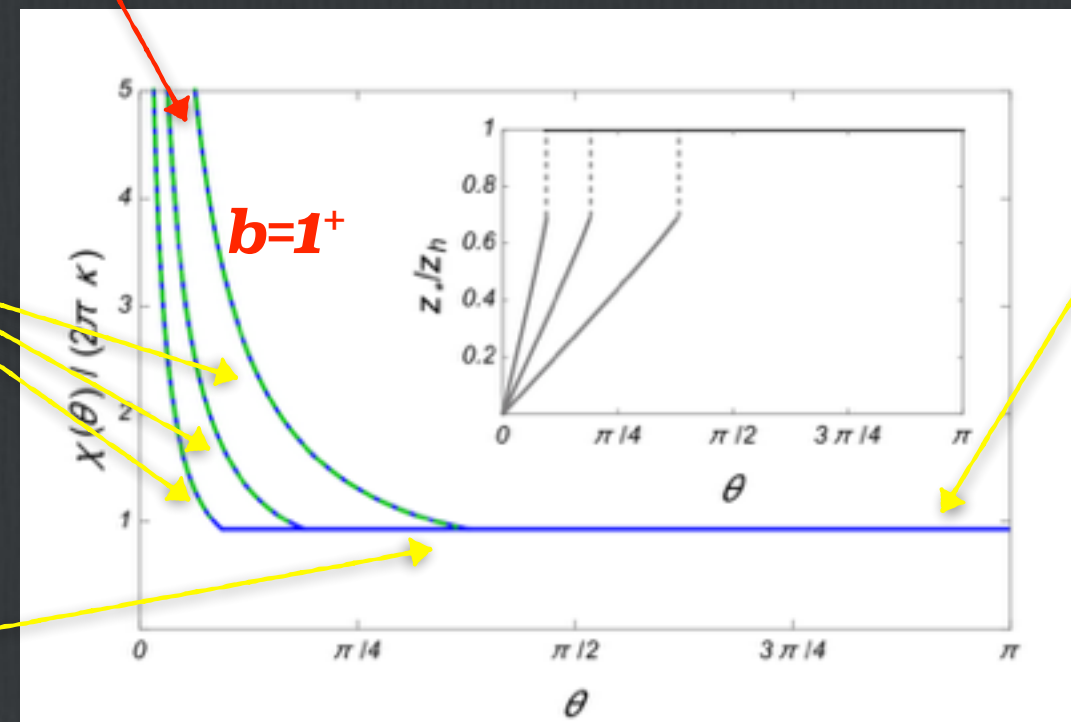
$$L_x > L_y$$

$$\chi(\theta) = \left[\frac{2\pi\kappa}{b} \right] \frac{1}{\theta} + \left[\frac{\Gamma(\frac{1}{4})^4 b^2}{432\pi G} \right] \theta^2 + \dots$$

$$0 < \frac{\theta}{2\pi} \leq \frac{p}{b} \quad p \simeq 0.1889$$

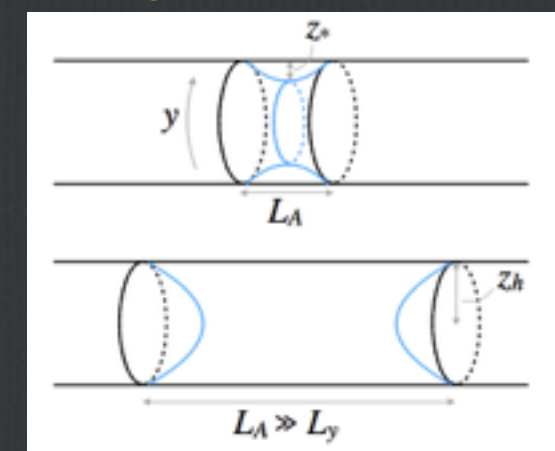
Non-smoothness (large-N)
Smooth for free scalar

Witczak-Krempa, Hayward, Melko



$$\chi(\theta) = \frac{2\pi}{3G}, \quad \frac{p}{b} < \frac{\theta}{2\pi} \leq \frac{1}{2}$$

Disconnected
holographic surface



2.2- Renormalized EE

Torus Entanglement

PB, Witczak-Krempa

Focus on cylinder limit with A covering half of it

$$S = \frac{L_y}{4G\epsilon} - \gamma$$

$$\gamma = \frac{\pi}{3G}$$

Pure constant

$$S_{\text{disk}} = \mathcal{B}R/\epsilon - \boxed{F}$$

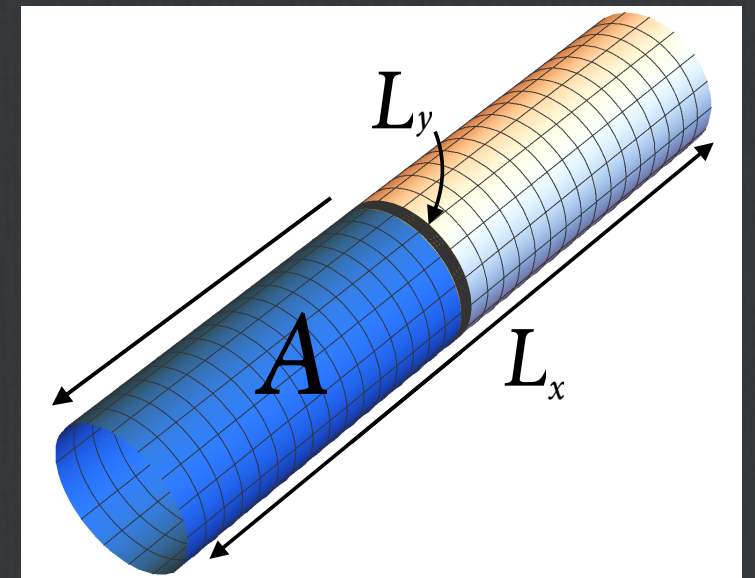


F-theorem

Casini, Huerta

Casini, Huerta, Myers, Yale

Very similar to disk entanglement



Monotonicity theorem for γ ?

Renormalized EE

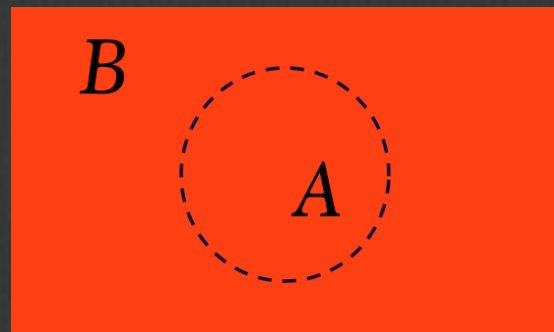
PB, Witczak-Krempa

Away from the fixed point, additional divergences

Solution \rightarrow use “Renormalized” EE

Liu, Mezei

Disk



$$S_{\text{disk}} = \mathcal{B}R/\epsilon - F$$

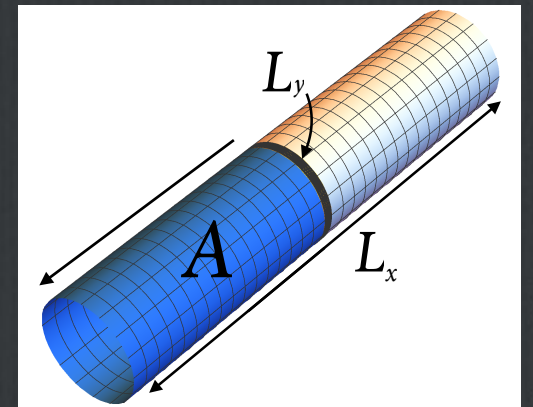


$$\mathcal{F}(R) = -S(R) + R \frac{\partial S(R)}{\partial R} \quad \text{Liu, Mezei}$$

Removes divergences
Isolates **F** at fixed point

This is in fact the quantity
that satisfies the **F-theorem**

Half cylinder



$$S_{\text{half cyl.}} = \mathcal{B}L_y/\epsilon - \gamma$$



$$\gamma_r^{(\alpha)}(L_y) = -S + L_y \frac{\partial S}{\partial L_y} + \alpha L_y^2 \frac{\partial^2 S}{\partial L_y^2}$$

Removes divergences
Isolates **gamma** at fixed points

(d=4 case is more interesting)

2.3- Holographic RG flow

Holographic RG flow

Deform CFT with relevant scalar perturbation

$$S = S_{\text{CFT}} + \lambda \int d^d x O(x) \rightarrow I = \int \frac{d^4 x \sqrt{-g}}{16\pi G} \left[6 + R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{m^2}{2} \phi^2 + \dots \right] \quad m^2 = \Delta(\Delta - 3)$$

Expansion parameter $\rightarrow \lambda z_{h0}^{3-\Delta} \propto \lambda L_y^{3-\Delta}$ unitary bound $1/2 < \Delta < 3$ marginality

Perturbed metric

$$ds^2 = \frac{1}{z^2} \left[\frac{dz^2}{f \cdot g_1(z)} + dx^2 + f \cdot g_2(z) dy^2 - dt^2 \right]$$

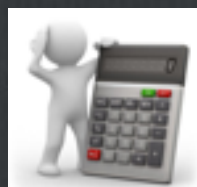
Regularity



$$z_h = \frac{3L_y}{4\pi} \sqrt{g_1(z_h)g_2(z_h)}$$

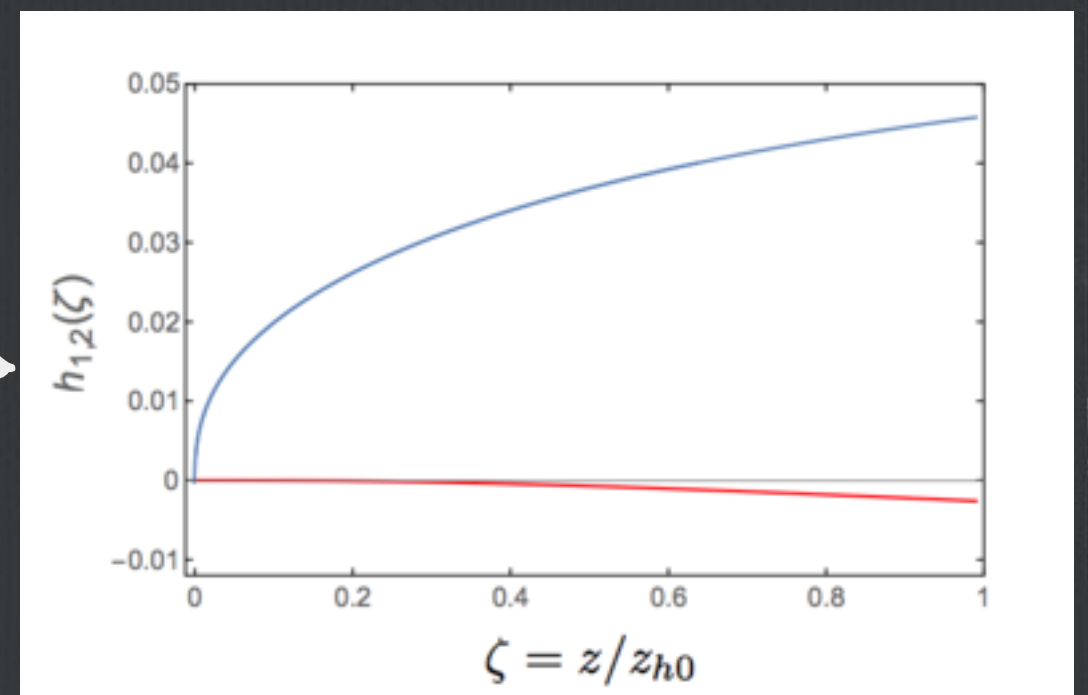
Scalar field back-reaction
at leading order fixes

do some
calculations!



$$g_i(z) = 1 + h_i(\zeta) z_{h0}^{2(3-\Delta)} \lambda^2 \rightarrow \text{where} \rightarrow$$

EE gets corrected in the perturbed geometry



Holographic RG flow

PB, Witczak-Krempa

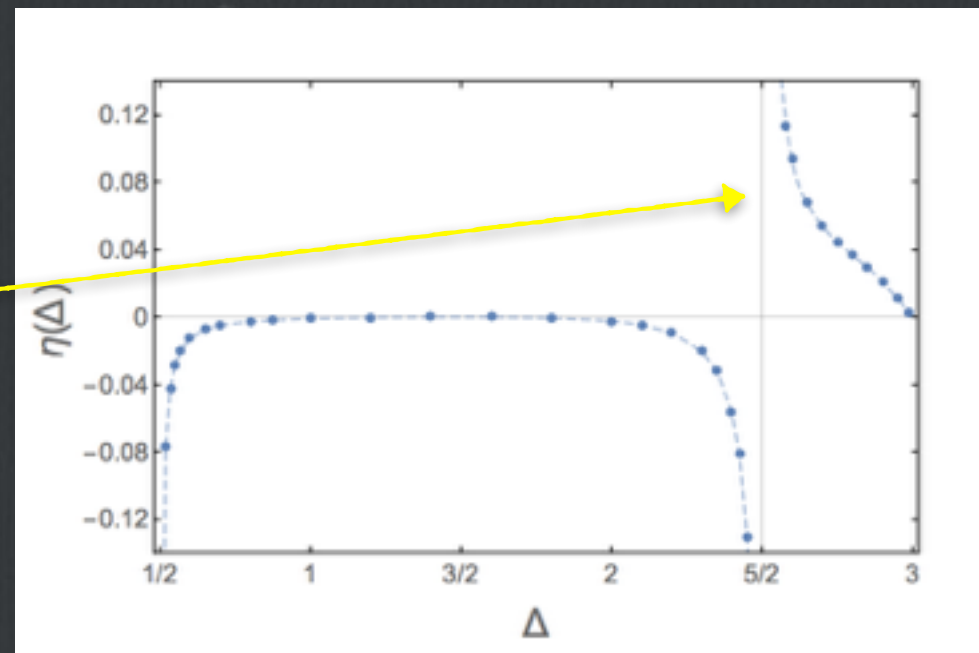
$$S = \frac{L_y}{2G\epsilon} - \frac{(3-\Delta)}{32(\Delta-5/2)} \frac{L_y \lambda^2}{\epsilon^{2\Delta-5} G} + \frac{(\frac{3}{4\pi})^{6-4\Delta} D_{0\Delta}}{4(2\Delta-1)} \frac{L_y^{7-4\Delta} \lambda^2}{\epsilon^{1-2\Delta} G} - 2\gamma + \dots$$

Extra divergences
Naive subtraction of
area law fails

Corrected universal
constant

$$\gamma = \frac{\pi}{3G} \left[1 - \eta(\Delta) L_y^{2(3-\Delta)} \lambda^2 \right]$$

Unphysical divergence
at $\Delta=5/2$



$$\eta(\Delta) = \left(\frac{3}{4\pi}\right)^{2(3-\Delta)} \left[\frac{h_1(1) + h_2(1)}{2} + \frac{3-\Delta}{16(\Delta-5/2)} - \frac{D_{0\Delta}}{2(2\Delta-1)} - \int_0^1 \frac{d\zeta}{2\zeta^2} \left(h_1(\zeta) - \frac{3-\Delta}{4} \zeta^{2(3-\Delta)} - D_{0\Delta} \zeta^{2\Delta} \right) \right]$$

Holographic RG flow

PB, Witczak-Krempa

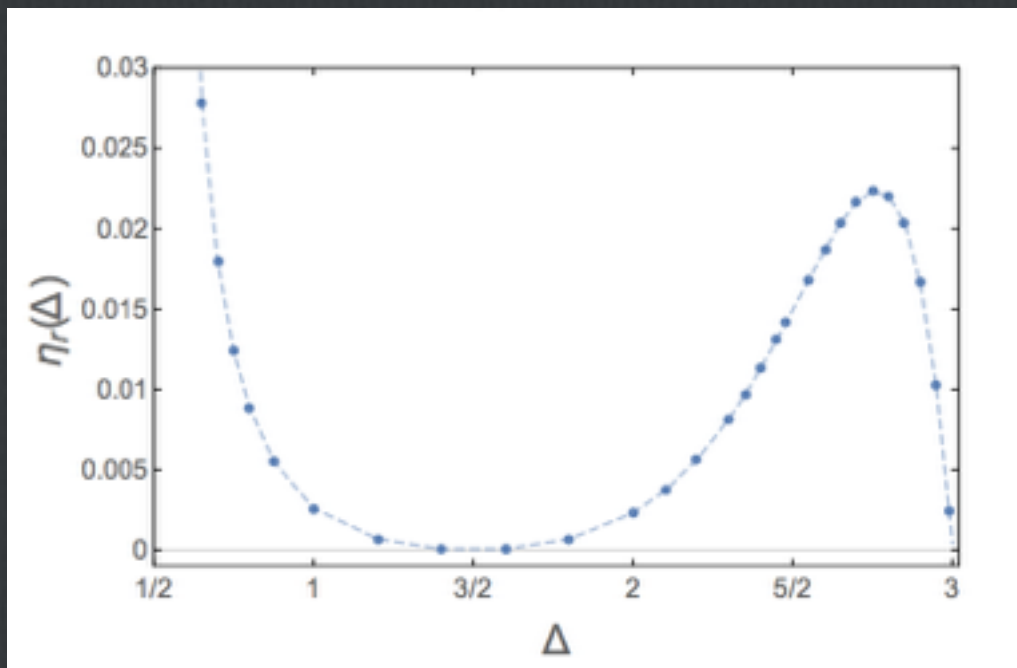
Use REE!

$$\gamma_r(L_y) = -S(L_y) + L_y \frac{\partial S(L_y)}{\partial L_y}$$

Removes all divergences
and issue at $\Delta=5/2$

$$\gamma_r(L_y) = \frac{\pi}{3G} \left[1 - \eta_r(\Delta) L_y^{2(3-\Delta)} \lambda^2 \right]$$

$$\eta_r(\Delta) = (2\Delta - 5)\eta(\Delta)$$



Decreasing for all Δ !

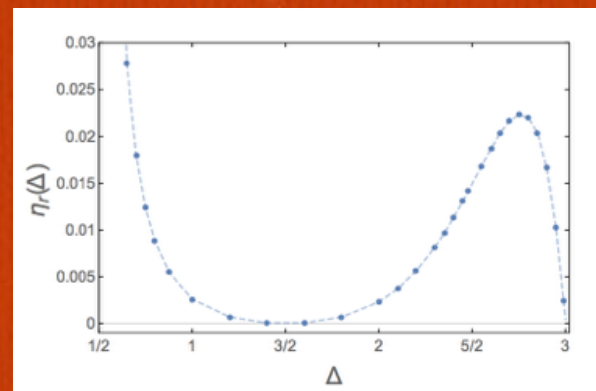
3- Summary

Defining good measures of degrees of freedom in CFTs is challenging

F-theorem for disk REE achieves it. But smooth regions not nice for simulations

Torus EE quite unexplored. Alternative to smooth regions.

In $d=3$ always decreasing at leading order for a particular holographic RG flow.



$$\gamma_r(L_y) = \frac{\pi}{3G} \left[1 - \eta_r(\Delta) L_y^{2(3-\Delta)} \lambda^2 \right]$$

What happens for more complicated flows and outside holography?

REE not uniquely defined in $d=4$.

$$\begin{aligned} \gamma_r^{(\alpha)}(L_1; r) = & -S(L_1) \\ & + (1 - \alpha) \frac{L_1}{2} \frac{\partial S(L_1)}{\partial L_1} + \alpha \frac{L_1^2}{2} \frac{\partial^2 S(L_1)}{\partial L_1^2} \end{aligned}$$



Thank you

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**Delta ITP visitor
programme**

d=4

$$S = \frac{L_1 L_2}{8G\epsilon^2} - \gamma$$

$$\gamma = \frac{\pi^2}{8G} \frac{L_2}{L_1}$$

Depends
on aspect ratio

Very different from smooth surface
entanglement in flat space

$$S_{\text{univ}} = I[a, c] \log(l/\epsilon)$$

↓
a-theorem

Myers, Sinha

Komargodski, Schwimmer (no use of EE)

**Free en.
on S^4
when A
sphere**

**Smooth
(sphere)**

$$S_{\text{sphere}} = \mathcal{B}(R/\epsilon)^2 - 4a \log(R/\epsilon)$$

$$\mathcal{S}(R) = \frac{1}{2} \left[R^2 \frac{\partial^2 S(R)}{\partial R^2} - R \frac{\partial S(R)}{\partial R} \right]$$

Liu, Mezei

Removes divergences
Isolates **a** at fixed points

Half cylinder

Scalable region
fix

Removes divergences
Isolates **γ** at fixed points

$$r = L_2/L_1$$

$$\gamma_r^{(\alpha)}(L_1; r) = -S(L_1) + (1 - \alpha) \frac{L_1}{2} \frac{\partial S(L_1)}{\partial L_1} + \alpha \frac{L_1^2}{2} \frac{\partial^2 S(L_1)}{\partial L_1^2}$$

Family of REE