

2. From random walk to probability distributions

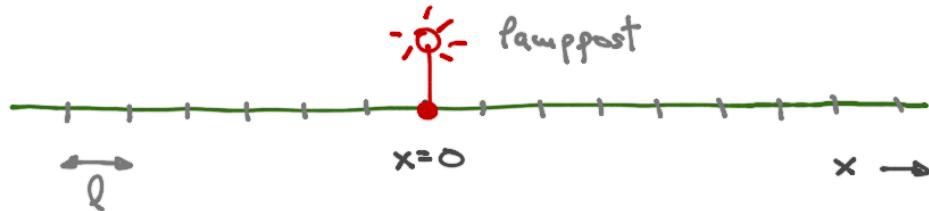
The aim of this chapter is to understand why the Gaussian probability distribution appears so often in Physics.

Statistical ensemble: It is an idealisation consisting of a large number of virtual copies of a system, considered all at once, each of which represents a possible state that the real system might be in. In other words, a statistical ensemble is a probability distribution for the state of the system. (ad)

I. A simple probabilistic system : Random Walk

The problem is stated as follows :

- * A drunk starts out from a lamppost ($x=0$) and each step is of equal length ℓ .
- * Step to right with probability p and to left with probability $q = 1-p$. Each step is independent of the preceding step.



- * After N steps the drunk will be at $x=m\ell$
- * Compute the probability (distribution) $P_N(m)$ of finding the drunk at position $x=m\ell$ after N steps.

Language: One can replace the drunk by a particle.

Denoting

$n_1 \equiv$ steps to the right

$n_2 \equiv$ steps to the left

one has

$$\begin{aligned} N &= n_1 + n_2 \\ m &= n_1 - n_2 \end{aligned} \quad \left\{ \begin{array}{l} n_1 = \frac{1}{2}(N+m) \\ n_2 = \frac{1}{2}(N-m) \end{array} \right.$$

Notice that

$$m = n_1 - n_2 = 2n_1 - N \Rightarrow \begin{cases} N \text{ odd} \Rightarrow m \text{ odd} \\ N \text{ even} \Rightarrow m \text{ even} \end{cases}$$

Probability of a sequence: Since there is no memory, the probability of any one given sequence of n_1 steps to the right and n_2 steps to the left is

$$\underbrace{p \cdots p}_{n_1} \underbrace{q \cdots q}_{n_2} = p^{n_1} q^{n_2}$$

But a given pair (n_1, n_2) can have multiple realisations so that the particle ends up at the same $x=m$ after the N steps

$$\text{Multiplicity} = \frac{N!}{n_1! n_2!}$$

Ex: $N=3$

				n_1	n_2	m	$\frac{N!}{n_1! n_2!}$
1)				3	0	3	1
2)				2	1	1	3
3)				1	2	-1	3
4)				0	3	-3	1
5)				1	2	-1	3
6)				2	1	1	3
7)				3	0	3	1
8)				2	1	1	3

The probability of taking (in a total of N step) n_1 steps to the right and $n_2 = N - n_1$ steps to the left is

$$W_N(n_1) = \underbrace{\frac{N!}{n_1! (N-n_1)!}}_{\text{multiplicity}} p^{n_1} (1-p)^{N-n_1}$$

Using $n_1 = \frac{1}{2}(N+m)$ yields

$$P_N(m) = W_N(n_1) = \frac{\frac{N!}{[\frac{1}{2}(N+m)]! [\frac{1}{2}(N-m)]!}}{p^{\frac{1}{2}(N+m)} (1-p)^{\frac{1}{2}(N-m)}}$$

Ex: In the case $N=3$ with $p=q=\frac{1}{2}$ one finds

$$P_3(m) = W_3(n_1) = \frac{1}{8}, \frac{3}{8}, \frac{3}{8}, \frac{1}{8}$$

The probability distribution above is called "binomial dist" by virtue of the binomial expansion

$$(p+q)^N = \sum_{n=0}^N \frac{N!}{n!(N-n)!} p^n q^{N-n}$$

II. Random walk and mean values

* Mean values in Mathematics : Let u be a variable that can take any of the M discrete values

$$u_1, u_2, \dots, u_M$$

with respective probabilities (stochastic variable)

$$P(u_1), P(u_2), \dots, P(u_M)$$

The mean or average value of u is defined by

$$\bar{u} = \frac{P(u_1)u_1 + P(u_2)u_2 + \dots + P(u_M)u_M}{P(u_1) + P(u_2) + \dots + P(u_M)} = \underbrace{\frac{\sum_{i=1}^M P(u_i)u_i}{\sum_{i=1}^M P(u_i)}}_{=1 \text{ as } P(u_i) \text{ is defined as a probability}}$$

$$= \sum_{i=1}^M P(u_i)u_i .$$

"normalisation cond"

More generically one has

$$\overline{f(u)} \equiv \sum_{i=1}^M P(u_i) f(u_i)$$

for a function $f(u)$.

• Property 1: $\overline{f(u) + g(u)} = \sum_{i=1}^n P(u_i) [f(u_i) + g(u_i)]$

$$= \sum_{i=1}^n P(u_i) f(u_i) + \sum_{i=1}^n P(u_i) g(u_i)$$

$$= \overline{f(u)} + \overline{g(u)}.$$

• Property 2: $\overline{cf(u)} = c \overline{f(u)}$ for $c = \text{cte.}$

There are some specific mean values which are particularly relevant in statistical physics

o) mean value $\bar{u} = \sum_{i=1}^n P(u_i) u_i$

1) first moment of u about its mean: $\overline{\Delta u} = \overline{(u - \bar{u})} = \bar{u} - \bar{u} = 0$

2) second moment of u about its mean or "dispersion"

$$\overline{(\Delta u)^2} \equiv \sum_{i=1}^n P(u_i) (u_i - \bar{u})^2 \geq 0 \quad \quad \quad \Delta u \equiv \sqrt{\overline{(\Delta u)^2}} \quad \text{= "Standard deviation"}$$

Note that

$$\overline{(\Delta u)^2} = \overline{(u - \bar{u})^2} = \overline{(u^2 - 2u\bar{u} + \bar{u}^2)} = \overline{u^2} - 2\bar{u}^2 + \bar{u}^2 = \overline{u^2} - \bar{u}^2$$

$$\Rightarrow \overline{(\Delta u)^2} = \overline{u^2} - \bar{u}^2 \geq 0 \Rightarrow \overline{u^2} \geq \bar{u}^2$$

n) n^{th} moment of u about its mean: $\overline{(\Delta u)^n}$

NOTE: Probability generating function $\Rightarrow g(t, u) = e^{tu}$

$$\Rightarrow \overline{u^k} = \frac{d^k g}{dt^k} \Big|_{t=0}$$

* Mean values for the random walk system : Let us focus on the binomial probability distribution

$$w(n_1) = \frac{N!}{n_1! (N-n_1)!} p^{n_1} q^{N-n_1}$$

- Normalisation :

$$\sum_{n_1=0}^N w(n_1) = \sum_{n_1=0}^N \frac{N!}{n_1! (N-n_1)!} p^{n_1} (1-p)^{N-n_1} = \underbrace{(p+q)^N}_{1} = 1$$

- o) Mean number of steps to the right :

$$\bar{n}_1 \equiv \sum_{n_1=0}^N w(n_1) n_1 = \sum_{n_1=0}^N \frac{N!}{n_1! (N-n_1)!} p^{n_1} q^{N-n_1} n_1 = (x)$$

Note : $n_1 p^{n_1} = p \frac{d}{dp} (p^{n_1})$

$$(x) = \sum_{n_1=0}^N \frac{N!}{n_1! (N-n_1)!} q^{N-n_1} p \frac{d}{dp} (p^{n_1})$$

$$= p \frac{d}{dp} \left[\sum_{n_1=0}^N \frac{N!}{n_1! (N-n_1)!} p^{n_1} q^{N-n_1} \right] = p \frac{d}{dp} [(p+q)^N]$$

$$= N p \underbrace{(p+q)^{N-1}}_1 = N p \Rightarrow \boxed{\bar{n}_1 = N p}$$

Note: Similarly one finds that $\bar{n}_2 = N q \Rightarrow \bar{n}_1 + \bar{n}_2 = N \underbrace{(p+q)}_1 = N$.

2) Dispersion :

$$\overline{(\Delta n_1)^2} = \overline{n_1^2} - \overline{n_1}^2 = \overline{n_1^2} - (N p)^2$$

This is what we have to compute

$$\overline{n_1^2} \equiv \sum_{n_1=0}^N w(n_1) n_1^2 = \sum_{n_1=0}^N \frac{N!}{n_1! (N-n_1)!} p^{n_1} q^{N-n_1} n_1^2 = (*)$$

Note: $n_1^2 p^{n_1} = \left[p \frac{d}{dp} \right]^2 (p^{n_1})$

$$\begin{aligned} (*) &= \left(p \frac{d}{dp} \right)^2 \sum_{n_1=0}^N \frac{N!}{n_1! (N-n_1)!} p^{n_1} q^{N-n_1} = \\ &= \left(p \frac{d}{dp} \right)^2 (p+q)^N = p \frac{d}{dp} \left(p^N (p+q)^{N-1} \right) \\ &= p^N \underbrace{(p+q)}_1^{N-1} + p^2 N(N-1) \underbrace{(p+q)}_1^{N-2} \underbrace{q}_{N-1} \\ &= p^N [1 + p(N-1)] = p^N [\underbrace{(1-p)}_{N-1} + p^N] \\ &= N p q + (N p)^2. \end{aligned}$$

So one finds

$$\overline{(\Delta n_1)^2} = \overline{n_1^2} - (N p)^2 = N p q \Rightarrow$$

$$\boxed{\overline{(\Delta n_1)^2} = N p q}$$

It is also meaningful to introduce the relative standard deviation

$$\frac{\Delta^* n_1}{\bar{n}_1} = \frac{\sqrt{\overline{(\Delta n_1)^2}}}{\bar{n}_1} = \frac{\sqrt{N p q}}{N p} = \sqrt{\frac{q}{p}} \cdot \frac{1}{\sqrt{N}} \xrightarrow{p=q} \frac{\Delta^* n_1}{\bar{n}_1} = \frac{1}{\sqrt{N}}$$

In terms of $m = n_1 - \bar{n}_1 = 2n_1 - N$ one has

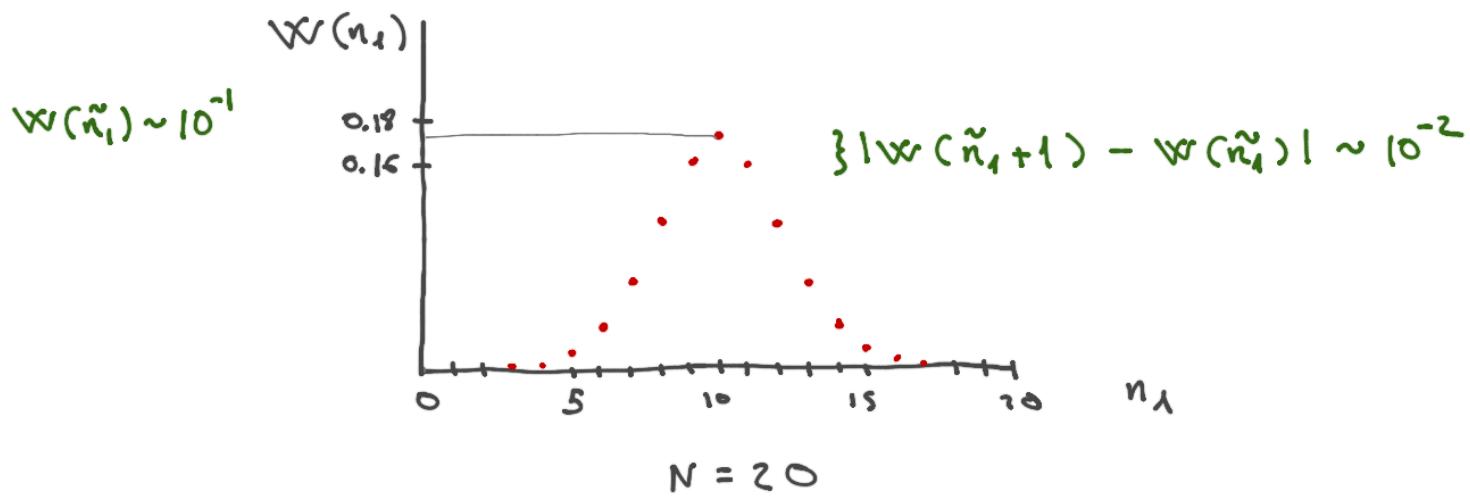
$$0) \bar{m} = 2\bar{n}_1 - N = 2Np - N = Np + N(p-1) = N(p-q)$$

$$1) \Delta m \equiv m - \bar{m} = (2n_1 - N) - (2\bar{n}_1 - N) = 2(n_1 - \bar{n}_1) = 2\Delta n_1 \Rightarrow (\Delta m)^2 = 4(\Delta n_1)^2$$

$$2) \overline{(\Delta m)^2} \equiv 4 \overline{(\Delta n_1)^2} = 4Npq$$

III. Probability distribution at large N

When $N \gg 1$ the probability distribution $W(n_1)$ has a sharp maximum around a value $n_1 = \tilde{n}_1$.



If N is large and one considers regions around the maximum $n_1 = \tilde{n}_1$, one has that

$$|W(n_1+1) - W(n_1)| \ll W(n_1)$$

Idea: W can, to a good approximation, be considered as a continuous function of the continuous variable n_1 although only $n_1 \in \mathbb{N}$ are of physical relevance.

The maximum of $W(n_1)$ is located at

$$\frac{dW}{dn_1} = 0 \quad \text{or equivalently} \quad \frac{d \ln W}{dn_1} = 0$$

Note: $\ln W$ is a much more slowly varying function.

Taylor expanding around $n_1 = \tilde{n}_1$ gives

$$\ln W(n_1) = \ln W(\tilde{n}_1) + B_1 (n_1 - \tilde{n}_1) + \frac{1}{2!} B_2 (n_1 - \tilde{n}_1)^2 + \dots$$

with

$$B_k = \left. \frac{d^k \ln W}{dn_1^k} \right|_{n_1 = \tilde{n}_1}$$

Since $n_1 = \tilde{n}_1$ is a maximum then $B_1 = 0$ and $B_2 < 0$ so

$$\ln W = \ln W(\tilde{n}_1) - \frac{1}{2} |B_2| (n_1 - \tilde{n}_1)^2 + \text{higher-order terms}$$

To first approximation one finds

$$W = \underbrace{\tilde{W}}_{W(\tilde{n}_1)} \cdot e^{-\frac{|B_2|}{2} (n_1 - \tilde{n}_1)^2}$$

Math: Starting from $W(n) = \frac{N!}{n! (N-n)!} p^n q^{N-n}$ then

$$\ln W = \ln N! - \log n! - \ln (N-n)! + n \ln p + (N-n) \ln q$$

Trick: If n is a large integer $n \gg 1$ then $\ln n!$ can be considered an almost continuous function of n since $\ln n!$ changes by a small fraction of itself if n is changed by a small integer. Hence

$$\frac{d \ln n!}{dn} \underset{\approx}{=} \frac{\ln(n+1)! - \ln n!}{1} = \ln \frac{(n+1)!}{n!} = \ln(n+1)$$

and for $n \gg 1$

$$\frac{d \ln n!}{dn} \underset{\approx}{=} \ln n$$

Then one has that

$$\frac{d}{dn_1} \ln w = -\ln n_1 + \ln(N-n_1) + \ln p - \ln q = \ln \left[\left(\frac{N-n_1}{n_1} \right) \frac{p}{q} \right]$$

Note 1: $\frac{d \ln (\tilde{n}-\tilde{n}_1)!}{d \tilde{n}_1} = -\frac{d \ln Y!}{d Y} \underset{\approx}{=} -\ln Y = -\ln(N-\tilde{n}_1)$

So the maximum is located at

$$\frac{d}{dn_1} \ln w \Big|_{n_1=\tilde{n}_1} = 0 \Rightarrow \frac{N-\tilde{n}_1}{\tilde{n}_1} = \frac{q}{p}$$

$$\Rightarrow (N-\tilde{n}_1)p = \tilde{n}_1 q$$

$$\Rightarrow Np = \tilde{n}_1 \underbrace{(p+q)}_{1}$$

$$\Rightarrow \tilde{n}_1 = Np = \bar{n}_1 \Rightarrow \begin{array}{l} \text{Maximum at} \\ \text{the mean value} \end{array}$$

Taking a further derivative

$$\frac{d^2}{dn_1} \ln W = \frac{n_1}{n-n_1} \frac{q}{p} \left(-\frac{N}{n_1^2} \frac{p}{q} \right) = -\frac{N}{n_1(n-n_1)}$$

and evaluating it at the maximum one gets

$$\begin{aligned} \left. \frac{d^2}{dn_1^2} \ln W \right|_{n_1=\tilde{n}_1} &= -\frac{N}{\tilde{n}_1(n-\tilde{n}_1)} = -\frac{N}{np(n-np)} \\ &= -\frac{N}{n^2 p(1-p)} = -\frac{1}{npq} \\ \Rightarrow B_2 &= -\frac{1}{npq} < 0 \end{aligned}$$

Finally the factor $\tilde{W} = W(\tilde{n}_1)$ can be determined from the normalisation condition

$$\sum_{n_1=0}^N W(n_1) \approx \int W(n_1) dn_1 = \int_{-\infty}^{\infty} W(\underbrace{\tilde{n}_1 + \eta}_{n_1}) d\eta \stackrel{!}{=} 1$$

NOTE: The integration over $-\infty < \eta < \infty$ is an excellent approximation since the contribution far from $\eta=0$ is negligible.

$$\int_{-\infty}^{\infty} W(n_1) d\eta \approx \int_{-\infty}^{\infty} \tilde{W} e^{-\frac{|B_2|}{2} \eta^2} d\eta = \tilde{W} \sqrt{\frac{2\pi}{|B_2|}} \stackrel{!}{=} 1$$

↳ see Math

$$\Rightarrow \tilde{\omega} = \sqrt{\frac{|B_2|}{2\pi}} = (2\pi N p_f)^{-1/2}$$

Math: Relevant integrals :

- $\int_0^\infty e^{-\eta^2} \eta^p d\eta = \Gamma(p+1) = p!$
- $\int_{-\infty}^\infty e^{-\eta^2} d\eta = \sqrt{\pi}$
- $I(p) \equiv \int_0^\infty e^{-\alpha\eta^2} \eta^p d\eta = \frac{1}{2} \Gamma\left(\frac{p+1}{2}\right) \alpha^{-\frac{p+1}{2}}$

so that

$$I(0) = \frac{1}{2} \sqrt{\pi} \alpha^{-1/2}$$

$$I(1) = \frac{1}{2} \alpha^{-1}$$

$$I(2) = \frac{1}{4} \sqrt{\pi} \alpha^{-3/2}$$

$$I(3) = \frac{1}{8} \alpha^{-2}$$

.....

Using these formulae we obtain

$$\int_{-\infty}^\infty e^{-\frac{1}{2}|B_2|\eta^2} d\eta = 2 \times I(0) = \sqrt{\pi} \frac{\sqrt{2}}{\sqrt{|B_2|}} = \sqrt{\frac{2\pi}{|B_2|}}$$

\downarrow
 $\alpha = \frac{|B_2|}{2}$

The previous computations show that

$$w(n_1) = (2\pi Npq)^{-1/2} e^{-\frac{(n_1-Np)^2}{2Npq}}$$

Recalling that $\bar{n}_1 = \tilde{n}_1 = Np$ and $\overline{(\Delta n_1)^2} = Npq$ then

$$w(n_1) = \left(2\pi \frac{\overline{(\Delta n_1)^2}}{Npq}\right)^{-1/2} e^{-\frac{(n_1-\bar{n}_1)^2}{2\overline{(\Delta n_1)^2}}}$$

"Gaussian distribution"

$(N \gg 1, n_1 \gg 1)$

IV. Gaussian probability distribution

Let us move to the more physical variable

$$x \equiv m \cdot l = (2n_1 - N)l \Rightarrow dx = l dm = 2l dn_1$$

$$\hookrightarrow m = n_1 - n_2$$

so that

$$w(x) = [2\pi Npq]^{-1/2} e^{-\frac{[\frac{x}{l} - N(p-q)]^2}{8Npq}}$$

Note: $(n_1 - Np) = \frac{1}{2}(N + m - 2Np) = \frac{1}{2}[m - N \underbrace{(2p-1)}_{(p-q)}] = \frac{1}{2}[m - N(p-q)]$

Equating the cumulative probability $P(x) dx = W(n_x) dn_x = W(x) \frac{dx}{2\ell}$

$$\begin{aligned}
 P(x) &= \frac{1}{2\ell} W(x) = \frac{1}{2\ell} [2\pi Npq]^{-1/2} e^{-\frac{[x-\mu-(p-q)\ell]^2}{8Npq}} \\
 &= \frac{1}{2\ell} [2\pi Npq]^{-1/2} e^{-\frac{[x-\mu-(p-q)\ell]^2}{\ell^2 8Npq}} \\
 &= \frac{1}{\sqrt{2\pi} \sigma} e^{-\frac{(x-\mu)^2}{2\sigma^2}} = P(m) \frac{dx}{2\ell}
 \end{aligned}$$

with $\mu \equiv (p-q)N\ell$ and $\sigma \equiv 2\sqrt{Npq}\ell$.

* Normalisation :

$$\int_{-\infty}^{\infty} P(x) dx = \frac{1}{\sqrt{2\pi} \sigma} \int_{-\infty}^{\infty} e^{-\frac{(x-\mu)^2}{2\sigma^2}} dx$$

$$\begin{aligned}
 y &= x-\mu & u &= \frac{1}{\sqrt{2\pi} \sigma} \int_{-\infty}^{\infty} e^{-\frac{y^2}{2\sigma^2}} dy \\
 dy &= dx & u &= \int_{-\infty}^{\infty} e^{-\frac{y^2}{2\sigma^2}} dy
 \end{aligned}$$

$$2\pi I(0) u = \frac{1}{\sqrt{2\pi} \sigma} 2 \frac{1}{2} \sqrt{\pi} \sqrt{2} \sigma = 1 \quad \checkmark$$

* Mean value of x :

$$\bar{x} = \int_{-\infty}^{\infty} P(x) x dx = \frac{1}{\sqrt{2\pi} \sigma} \int_{-\infty}^{\infty} x e^{-\frac{(x-\mu)^2}{2\sigma^2}} dx$$

$$\begin{aligned}
 &= \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{\infty} (y+\mu) e^{-\frac{y^2}{2\sigma^2}} dy \\
 &= \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{\infty} y e^{-\frac{y^2}{2\sigma^2}} dy + \mu \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{\infty} e^{-\frac{y^2}{2\sigma^2}} dy \\
 &\quad \underbrace{\qquad\qquad\qquad}_{\text{zero} \equiv \text{odd function}} \qquad\qquad\qquad \underbrace{\qquad\qquad\qquad}_{2 \times I(0)}
 \end{aligned}$$

$$= \mu \frac{1}{\sqrt{2\pi}\sigma} + \frac{1}{2} \sqrt{\pi} \sqrt{2}\sigma = \mu = (p-q)Nl$$

* Dispersion :

$$\begin{aligned}
 \overline{(\Delta x)^2} &= \overline{(x-\mu)^2} = \int_{-\infty}^{\infty} (x-\mu)^2 P(x) dx \\
 &= \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{\infty} y^2 e^{-\frac{y^2}{2\sigma^2}} dy \\
 2 \times I(2) \quad \omega &= \frac{1}{\sqrt{2\pi}\sigma} \underbrace{\frac{1}{2}}_{\frac{1}{2}} \sqrt{\pi} \underbrace{(2\sigma^2)^{3/2}}_{\sqrt{2}\sigma^2\sigma^2} = \sigma^2 = 4Npq l^2
 \end{aligned}$$

TAKE HOME MESSAGE : Starting from the simplest probabilistic system (random walk) and taking $N \gg 1$ and $n \gg 1$ limits we found that a Gaussian distribution emerges !!

v. Poisson probability distribution

It appears also as a limit of the binomial distribution

$$N \gg 1, p \ll 1 \text{ so } n_1 \ll N$$

Starting from

$$W(n_1) = \frac{N!}{n_1! (N-n_1)!} p^{n_1} q^{N-n_1} \underset{\substack{\text{underbrace} \\ \text{underbrace}}}{\approx} \frac{N^{n_1}}{n_1!} p^{n_1} e^{-Np} = \underbrace{\frac{\lambda^{n_1}}{n_1!} e^{-\lambda}}$$

$$\frac{N(N-1)(N-2)\cdots(N-n_1+1)}{n_1!} \underset{\substack{\downarrow \\ N \gg n_1 \\ N \gg 1}}{\approx} \frac{N^{n_1}}{n_1!}$$

with $\lambda \equiv Np$.

NOTE: $\log q^{N-n_1} = (N-n_1) \log q = (N-n_1) \log(1-p)$

$$= N(-p + \text{higher order}) = -Np$$

$p \ll 1 \rightarrow$
 $N \gg n_1$
 $\Rightarrow q^{N-n_1} = e^{-Np}$

NOTE: $\log(1+x) = x - \frac{x^2}{2} + \frac{x^3}{3} + \dots$

* Normalisation : $\sum_{n_1=0}^N \frac{\lambda^{n_1}}{n_1!} e^{-\lambda} \underset{N \gg 1}{=} e^{-\lambda} \sum_{n_1=0}^{\infty} \frac{\lambda^{n_1}}{n_1!} = e^{-\lambda} e^{\lambda} = 1 \quad \checkmark$

* Mean value of n_1 :

$$\bar{n}_1 = \sum_{n_1=0}^N n_1 \frac{\lambda^{n_1}}{n_1!} e^{-\lambda} = \sum_{n_1=1}^N n_1 \frac{\lambda^{n_1}}{n_1!} e^{-\lambda}$$

$$= e^{-\lambda} \lambda \sum_{n_1=1}^N \frac{\lambda^{n_1-1}}{(n_1-1)!} = e^{-\lambda} \lambda \sum_{m=0}^{\infty} \frac{\lambda^m}{m!} = \lambda$$

$\underbrace{e^{-\lambda}}$

$n_1-1=m$
 $N \gg 1$

* Dispersion:

$$\overline{(\Delta n_1)^2} = \overline{n_1^2} - \bar{n}_1^2 = \lambda(\lambda+1) - \lambda^2 = \lambda$$

$\underbrace{\quad}_{(4)}$

$$(4) \overline{n_1^2} = \sum_{n_1=0}^N n_1^2 \frac{\lambda^{n_1}}{n_1!} e^{-\lambda} = e^{-\lambda} \sum_{n_1=1}^N n_1 \frac{\lambda^{n_1-1}}{(n_1-1)!}$$

$$= e^{-\lambda} \lambda \frac{d}{d\lambda} \sum_{n_1=1}^N \frac{1}{(n_1-1)!} \lambda^{n_1-1} \cdot \lambda = e^{-\lambda} \lambda \frac{d}{d\lambda} (e^\lambda \cdot \lambda)$$

$n_1-1=m$
 $N \gg 1$

NOTE: $\lambda \frac{d \lambda^{n_1}}{d\lambda} = n_1 \lambda^{n_1}$

$$= e^{-\lambda} \lambda (e^\lambda \lambda + e^\lambda) = \lambda(\lambda+1).$$

TAKE HOME MESSAGE: The Poisson distribution has to do with very unlikely (rare) events that happen to occur because of a very large population ($N \gg 1$).