De Sitter Universes

from

String Theory

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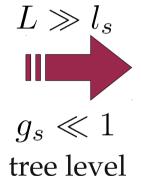
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Linking strings to the world

- String Theory = best known framework where to describe GR + QFT
- > The fundamental building blocks are tiny strings with $l_s \sim 10^{-33} {
 m cm}$

low energies

String Theory



Classical Field Theory

- ten dimensional space-time
- supersymmetry

$$(G_{MN}, A_M, \varphi)$$
 + fermi

4D world $\Lambda_{c.c} > 0$



10D SUPERGRAVITY!!

Six extra dimensions!!

The footprint of the extra dimensions

Fluctuations of the internal space around a fixed geometry translates into massless 4d scalar fields known as ``moduli''

$$\mathcal{L} = \frac{1}{2} R - \frac{1}{2} \partial_{\mu} \phi_i \, \partial^{\mu} \phi^i$$



massless scalars = long range interactions (precision tests of GR)

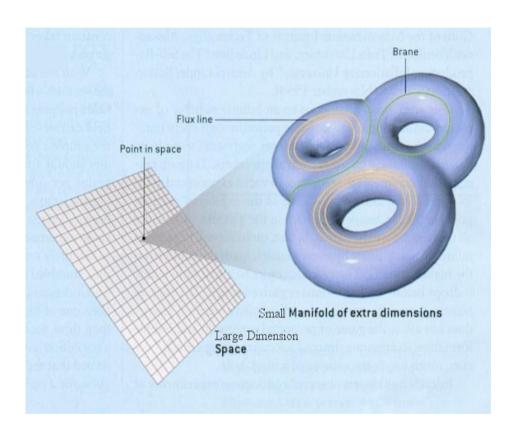
$$V(\phi) = m_{ij}^2 \, \phi^i \phi^j + \dots$$

$$m_{ij}^2 > 0$$

Moduli VEVs $\langle \phi \rangle = \phi_0$ determine 4d physics !!

$$\Lambda_{c.c} \equiv V(\phi_0)$$
 $\left(g_s, \operatorname{Vol}_{int}\right)$ fermi masses

Extra dimensions...



... will be non empty!!

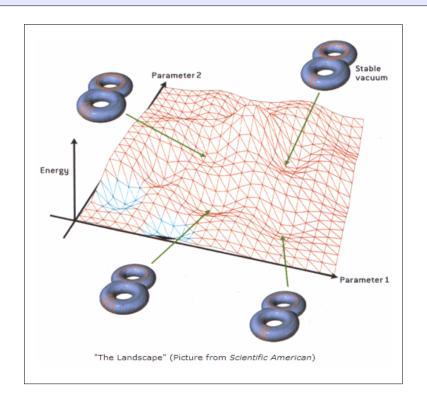
- > D-branes $\left(M \propto 1/g_s\right)$
- magnetic fluxes
- funny geometries

• •

$$V(\phi) = V_{brane} + V_{flux} + V_{geom} + \dots$$

The problem

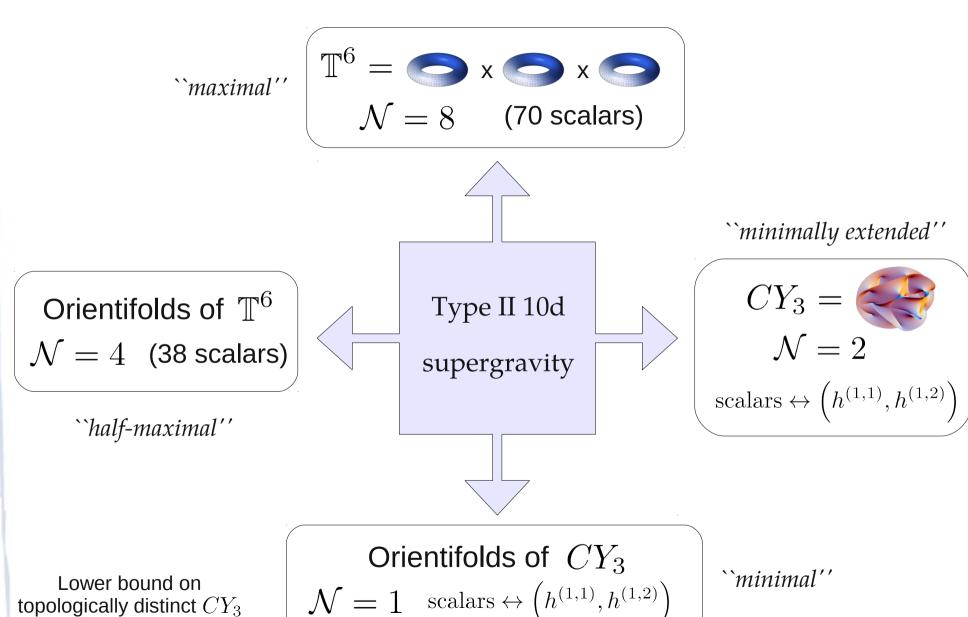
Model building:



$$\Lambda_{c.c} \equiv V(\phi_0) > 0$$

... but where is de Sitter within the string landscape ??

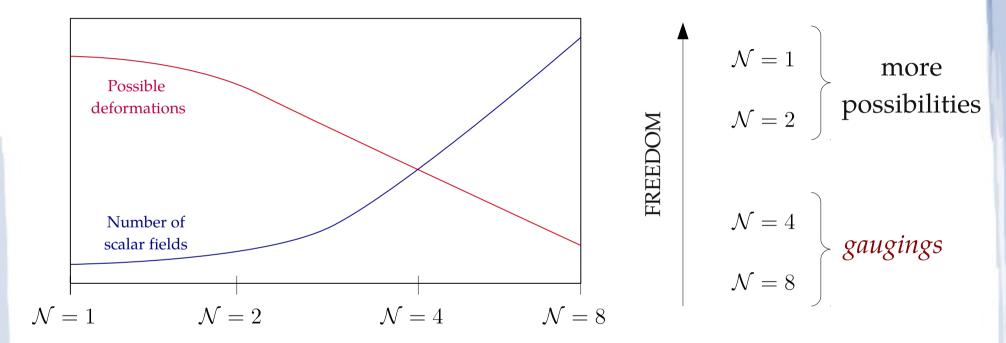
Internal geometries and massless theories . . .



topologically distinct CY_3 30108

... how to deform massless theories to have $V(\phi) \neq 0$?

Supersymmetry dictates what deformations are allowed



gaugings = part of the global symmetry is promoted to local (gauge)

Questions:

- > Which deformations induce potentials with stable De Sitter extrema?
- Where do these deformations come from in string theory ?

De Sitter in extended supergravity

> $\mathcal{N}=8$: unstable dS solutions with SO(4,4) and SO(5,3) gaugings [Hull, Warner '85]

 $\mathcal{N}=4$: unstable dS solutions with gaugings at angles

[De Roo, Wagemans '85]

$$i$$
) $G_1 imes G_2$ gaugings with $\begin{cases} G_i = SO(p_i,q_i) \;\;,\;\; p_i+q_i=4 \ \\ G_i = CSO(p_i,q_i,r_i) \;\;,\;\; p_i+q_i+r_i=4 \end{cases}$

ii) $SO(3,1) \ltimes U(1)^6$ gauging

[De Roo, Westra, Panda, (Trigiante) '02, '03, '06]

[Dibitetto, A.G, Roest '11]

non-geometric fluxes in string theory!!

[Dibitetto, Linares, Roest '10]

> $\mathcal{N}=2$: stable dS solutions with $SO(2,1)\times SO(3)$ gauging plus Fayet-Iliopoulos terms [Fré, Trigiante, Van Proeyen '03]

unclear origin in string theory!!

De Sitter in minimal supergravity

* No-go theorems forbidding dS solutions in $\mathcal{N}=1$ compactifications with magnetic fluxes

 $V_o = -\frac{1}{9} \sum \bar{F}^2 \le 0$ AdS!!

[Hertzberg, Kachru, Taylor, Tegmark '07]

Fincluding more general fluxes : (metric + non-geometric)

$$V_o = -\frac{1}{9} \sum \bar{F}^2 + \Delta V_{\text{metric}} + \Delta V_{\text{non-geom}}$$

a) metric fluxes **\ unstable** dS in type IIA models

[Caviezel, Koerber, Kors, Lust, Wrase, Zagerman '08]

b) non-geometric fluxes \longleftrightarrow stable dS in type IIA models

[de Carlos, A.G, Moreno '09, '10]

> Including D-branes to uplift an AdS solution

[Kachru, Kallosh, Linde, Trivedi '03]

a) D-terms from D-branes \iff stable dS in type IIB models

[Burgess, Kallosh, Quevedo '03]

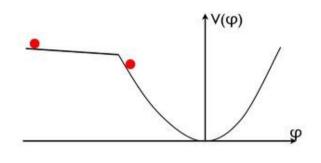
b) non-perturbative effects from D-branes \iff stable dS in type IIB

[Achúcarro, de Carlos, Casas, Doplicher '06]

Cosmology from moduli?

' slow-roll inflation requires an almost flat dS saddle point of $\,V(\phi)\,$ from which to start rolling down

$$\eta \equiv M_p^2 \left(\frac{V''}{V}\right) \ll 1$$



- $^{\flat}$ dS saddle points suffering from eta-problem, *i.e.* $\eta \sim \mathcal{O}(1)$
 - *i*) gaugings in extended supergravity

[Kallosh, Linde, Prokushkin, Shmakova '01]

ii) general fluxes in minimal supergravity

[Flauger, Paban, Robbings, Wrase '08] [de Carlos, A.G, Moreno '10]

 $^{\diamond}$ dS saddle points with $\eta \ll 1$ in minimal supergravity including non-perturbative effects \implies axion inflation!!

[Dimopoulos, Kachru, McGreevy, Wacker '05]

Overview

- > Finding dS solutions from string compactifications represents a crucial step in connecting strings with real world physics
- $^{\flat}$ All known dS solutions coming from gaugings in $~\mathcal{N} \geq 4~$ turn out to be unstable. They correspond to string compactifications in the presence of non-geometric fluxes for which a higher-dimensional origin is still unclear
- > Some stable dS solutions are found in $\mathcal{N} \leq 2$ considering gaugings together with additional ingredients. Their higher-dimensional origin is not well understood either
- [>] D-branes seem to be crucial in order to find/understand stable dS solutions
- [>] Cosmology certainly represents an interesing playground to test string theory

Thanks for your attention!!

Extra dimensions: compact vs non-compact

compactification

$$ds_{(5)}^2 = ds_{(4)}^2 + dy^2$$

$$M_{(4)}^2 = M_{(5)}^3 L_y$$

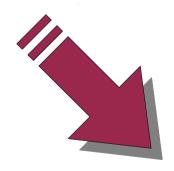
[Kaluza '21, Klein 26']

warped non-compact

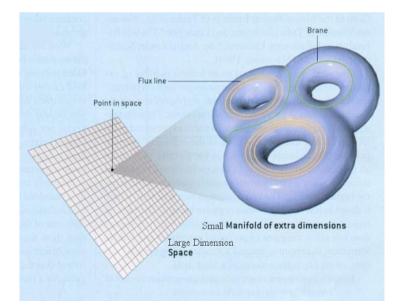
$$ds_{(5)}^2 = e^{-2k|y|} ds_{(4)}^2 + dy^2$$

$$M_{(4)}^2 = \frac{M_{(5)}^3}{k} \left(1 - e^{-2kL_y}\right)$$

[Randall, Sundrum '99]



10d space-time



- ... will be non empty!!
- funny geometries
- * D-branes $\left(M \propto 1/g_s\right)$
- Magnetic fields

Gaugings and their higher-dimensional origin

- Scalars potentials are induced by ``gaugings'': Part of the global symmetry is promoted to local (gauge)

 [de Wit, Samtleben, Trigiante '07]
 [Schon, Weidner '06]
- $\mathcal{N}=8$: Gauging a subgroup of the global symmetry $G=E_7$

Internal space extension



Exceptional Generalised Geometry?

[Pacheco, Waldram '08, Grana, Louis, Sim, Waldram '09] [Aldazabal, Andrés, Cámara, Grana '10]

 $^{\triangleright} \mathcal{N} = 4$: Gauging a subgroup of the global symmetry $\mathit{G} = \mathit{SL}(2) \times \mathit{SO}(6,6)$

Internal space extension



Doubled/Generalised Geometry?

[Hitchin '02, Gualtieri '04] [Hull '04, '06]



String compactifications including generalised flux backgrounds !!